

Photoproduction of η mesons from the neutron: cross sections and double polarization observable E

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Abstract. Results from measurements of the photoproduction of η mesons from quasifree protons and neutrons are summarized. The experiments were performed with the CBELSA/TAPS detector at the electron accelerator ELSA in Bonn using the $\eta \rightarrow 3\pi^0 \rightarrow 6\gamma$ decay. A liquid deuterium target was used for the measurement of total cross sections and angular distributions. The results confirm earlier measurements from Bonn and the MAMI facility in Mainz about the existence of a narrow structure in the excitation function of $\gamma n \rightarrow n\eta$. The current angular distributions show a forward-backward asymmetry, which was previously not seen, but was predicted by model calculations including an additional narrow P_{11} state. Furthermore, data obtained with a longitudinally polarized, deuterated butanol target and a circularly polarized photon beam were analyzed to determine the double polarization observable E . Both data sets together were also used to extract the helicity dependent cross sections $\sigma_{1/2}$ and $\sigma_{3/2}$. The narrow structure in the excitation function of $\gamma n \rightarrow n\eta$ appears associated with the helicity-1/2 component of the reaction.

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1 Introduction

The excitation spectrum of nucleons is one of the most important testing grounds of our understanding of the strong interaction in the non-perturbative regime [1,2]. While in the past, most experimental information came from hadron induced reactions such as elastic and inelastic pion scattering (profiting from the large reaction cross sections), the last two decades have seen a huge worldwide effort to study the electromagnetic excitations of

baryons with photon induced reactions [3,4] and with electron scattering [5,6].

This program has two central parts. The study of many different final states should eliminate experimental bias and the measurement of several types of polarization observables in addition to differential cross sections should provide data sets that allow for (almost) uniquely determined partial wave analyses. Since the electromagnetic excitation of nucleon resonances depends on isospin, it is not sufficient to study only photon induced reactions on protons, but the neutron target must also be investigated. This part of the experimental program is still much less advanced than the measurements with free proton targets.

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It involves experiments with quasifree neutrons (mostly bound in the deuteron) with the difficulties from the co-incident detection of the recoil nucleons, Fermi motion of the bound nucleons, and effects from final state interactions (FSI), complicating the interpretation of the results. However, recent progress was encouraging (see e.g. [7,8]), and currently such experiments are in the focus at the Bonn ELSA and Mainz MAMI facilities.

Photoproduction of η mesons is a selective reaction and due to the isoscalar nature of this meson, only N^* resonances can decay to the nucleon ground state via η emission (higher lying Δ states can emit η mesons in decays to $\Delta(1232)\eta$ resulting in a $N\eta\pi$ final state [9]). The main features of the $\gamma p \rightarrow p\eta$ reaction at low incident photon energies (threshold is at $E_\gamma \approx 707$ MeV) are well understood. This reaction is dominated by the excitation of the $N(1535)1/2^-$ and $N(1650)1/2^-$ states [10,11] and a very small contribution from the $N(1520)3/2^-$ state has been identified mainly from the $S_{11} - D_{13}$ interference term in the photon beam asymmetry [12,13,14]. Angular distributions for $\gamma p \rightarrow p\eta$ from threshold up to photon energies of 2.5 GeV were measured by several experiments, in particular at CLAS [15,16], ELSA [17,18,19], GRAAL [20], LNS [21], and at MAMI [10,22]. Recently, also results for the transverse target asymmetry T and the beam-target asymmetry F were published from a MAMI experiment [23] and results for the double polarization observable E became available from CLAS [24]. The status of η production from nucleons and nuclei with hadron and photon induced reactions was recently reviewed in [25].

Photoproduction of η mesons off neutrons is an intriguing example of why measurements with neutron targets are important. Measurements with deuterium and helium targets in the threshold region [26,27,28,29,30,31,32,33] determined the isospin structure of the electromagnetic excitation of the dominant $N(1535)1/2^-$ resonance [3]. But a real surprise came from the first studies of the $\gamma n \rightarrow n\eta$ reaction at somewhat higher energies, which revealed an unexpected, narrow structure in the excitation function around incident photon energies of 1 GeV ($W \approx 1.67$ GeV). This structure was first reported by the GRAAL collaboration [34] and confirmed by the CBELSA/TAPS [35,36] collaboration at ELSA and also at LNS in Sendai [37]. Recent high-statistics measurements at MAMI using a deuterium and ^3He target [38,39,40] provided much more precise results for the excitation functions and the angular distributions. These measurements determined the position of the structure at a final state invariant mass of $W = (1670 \pm 5)$ MeV and a width of only $\Gamma = (30 \pm 15)$ MeV. The nature of this structure is not yet understood. Some analyses suggested a contribution from a new narrow nucleon resonance [41,42,43,44,45]. In its 2014 edition, the Review of Particle Physics (RPP) [46] lists this as a tentative $N(1685)$ state with unknown quantum numbers, but from the 2016 version [47] it was removed again. However, other models use coupled-channel effects of known nucleon resonances [48,49] or intermediate strangeness states [50]. A fit [51] to the high statistics A2 deuteron data [38,40] by the BnGa

group suggests an interference in the $J^P = 1/2^-$ S partial wave between contributions of the well-known $N(1535)$ and $N(1650)$ resonances [51]. This solution requires a sign change of the electromagnetic coupling of the $N(1650)$ for the neutron compared to the results listed in the RPP [47]. Alternative fits with a narrow P_{11} -like $N(1685)$ resonance were seen as inferior since they yielded a larger χ^2 [51].

In the meantime, the situation became even more complicated. A footprint of this narrow structure was also found in the beam asymmetry of Compton scattering from protons [52]. A further narrow structure at an invariant mass of $W \approx 1.726$ GeV in Compton scattering [52] was also identified in the $\gamma n \rightarrow n\eta$ reaction [53]. There are also pronounced structures in this energy range in the Legendre expansion coefficients of the beam asymmetry for η production off the proton [25]. The appearance of such structures in quite different reactions makes explanations with intricate interference effects less probable. It is obvious that the $\gamma n \rightarrow n\eta$ reaction deserves a lot of attention in this energy range.

The results as a function of the incident photon energy for the total cross sections and angular distributions measured at ELSA in Bonn [35,36] and at MAMI in Mainz [38,40] agreed. However, there was some discrepancy in the absolute scale of the cross sections when the effective W was reconstructed from the final state kinematics. It was already discussed in [25] that there is some internal inconsistency between the results as function of E_γ and as function of W (although just inside quoted systematic uncertainties) in [36]. The results as function of reconstructed W from the current work, measured with a very similar setup as the results from [36] but with a much finer binning in the angular distributions than in [36] supersede the previous ones from ELSA.

More information about the interesting structures in the $\gamma n \rightarrow n\eta$ reaction can be obtained with the measurement of polarization observables. So far, only the beam asymmetry Σ has been investigated with the GRAAL experiment [54]. The particularly interesting double polarization observable E can be measured with a circularly polarized photon beam and a longitudinally polarized target. It is defined as:

$$E = \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{\sigma_{1/2} - \sigma_{3/2}}{2\sigma_0}, \quad (1)$$

where $\sigma_{1/2}$ and $\sigma_{3/2}$ are the helicity dependent cross sections with photon and nucleon spin antiparallel or parallel, respectively, and σ_0 is the unpolarized cross section. The helicity dependent cross sections, which can be easily computed when E and σ_0 are known, are especially interesting, since they allow direct access to the spin structure. Resonances with spin $J = 1/2$ (like S_{11} and P_{11}) contribute only to $\sigma_{1/2}$. States with spin $J \geq 3/2$ can also appear in $\sigma_{3/2}$ and very likely will. There are no known examples of $J \geq 3/2$ states which contribute only to $\sigma_{1/2}$, most of them have larger contributions to $\sigma_{3/2}$ [47]. Very recently, results for the quasifree proton and neutron were published by the A2 collaboration at MAMI [55]. Here, we

will report results from a measurement of this observable at ELSA.

2 Experimental Setup

The measurements for the unpolarized cross sections and the double polarization observable E were carried out by the CBELSA/TAPS collaboration at the electron stretcher accelerator ELSA [56] in Bonn, Germany. The accelerator provided an electron beam with an energy of $E_0 = 2.35$ GeV.

The circularly polarized photons were generated from longitudinally polarized electrons via incoherent bremsstrahlung on a thin radiator foil (20 μm Vacoflux50). The radiator foil was magnetized and allowed for constant monitoring of the electron polarization during data taking with the help of Møller scattering. For the current experiments, the polarization of the electrons was in the range $P_e = 60\text{--}65\%$ and was determined with an uncertainty of 2% [57].

The circular photon polarization P_γ was deduced from the longitudinal electron polarization with the following polarization transfer equation [58]:

$$\frac{P_\gamma}{P_e} = \frac{3 + (1 - x)}{3 + 3(1 - x)^2 - 2(1 - x)} \cdot x, \quad (2)$$

where $x = E_\gamma/E_0$, and E_γ is the energy of the photon. The photon energies were determined by the tagging spectrometer, which is mainly a dipole magnet and analyzes the momentum of the scattered electrons. The tagger consists of 480 fibers of 2 mm diameter and 96 scintillating bars with a thickness of 1.4–5 cm. This setup enabled the energy tagging of the photons in the energy range from 0.442 GeV to 2.3 GeV with a resolution of 2–25 MeV [59]. Typical tagger rates were 6.7 MHz for the measurement of the unpolarized cross sections and 10.5 MHz for the measurement of the polarization observable E .

The energy tagged photons impinged on a target, which was centered in the Crystal Barrel (CBB) detector, as shown in Fig. 1. The target for the measurement of the unpolarized cross section was similar to earlier measurements with a liquid deuterium target [36]. The target cell was a capton cylinder (0.125 mm foil thickness) with a diameter of 3 cm and a length of 5.26 cm. It was filled with LD_2 with a density of 0.169 g/cm³ corresponding to a surface thickness of 0.26 nuclei/barn. For the double polarization experiment, a 1.88 cm long deuterated butanol ($\text{C}_4\text{D}_9\text{OD}$) target with an effective density of 0.65 g/cm³ was used. The deuterated butanol was polarized with Dynamic Nuclear Polarization (DNP) [60]. The target polarization was 50–60% for the current work and was measured with an NMR coil every two days with an overall relative uncertainty of 5%. The polarization values between the measurements were interpolated using an exponential function with a decay time of approximately 340 hours. Data from each single run of data taking (approximately 30 min) were normalized individually to the beam and target polarization degrees determined for this run.

The main part of the solid angle ($30^\circ \leq \theta \leq 156^\circ$ and $0^\circ \leq \phi \leq 360^\circ$) was covered by the CBB calorimeter [61], which consisted of 21 azimuthally symmetric rings of 60 CsI(Tl) crystals each. Hence, an angular coverage of $\Delta\phi = \Delta\theta = 6^\circ$ (for the 21st ring $\Delta\phi = 12^\circ$) was provided and an angular resolution of 1.5° was achieved by using the center of gravity of a shower. The energy resolution was given by $\sigma_E/E = 1.5/\sqrt[4]{E[\text{GeV}]}$ [61]. In the CBB, charged particles were detected using the inner detector [62], which was made of three layers of 513 scintillating fibers. By requiring at least two intersecting hit fibers, an angular resolution of $\Delta\theta = 0.4^\circ$ and $\Delta\phi = 0.1^\circ$ was reached.

Up to an angle of $\theta = 11.2^\circ$, the CBB was complemented with the Forward Plug (FP) detector made of the same CsI(Tl) crystals as the CBB (actually these are the modules of the old configuration of the CBB [61], which had been removed from the detector when it was used together with a larger forward wall from the TAPS detector as in [36]). In contrast to the CBB, the FP detector provided time information due to the photomultiplier (PMT) readout (CBB was read out by photodiodes). In front of the FP crystals, two layers of plastic scintillators were mounted to identify charged particles.

The most forward region ($2^\circ \leq \theta \leq 12^\circ$) of the experiment was covered by the MiniTAPS (MT) detector, which is a subunit of the TAPS detector [63,64] made of 216 BaF₂ crystals read out via PMTs. Each crystal had its own plastic veto to identify charged particles. A gas Cerenkov detector with $n = 1.00043$, mounted between the CBB and the MT detectors, was used as an online veto to reduce electromagnetic background in the trigger.

This detector arrangement covered almost the complete solid angle for the detection of photons and recoil nucleons, but not for triggering on them, because the missing time information from the CBB did not allow to use it in the first-level trigger. This was not a problem for reactions with recoil protons because in that case, the inner scintillating fiber detector can trigger on the protons. But for final states composed of photons and neutrons, only the FP and the MT could trigger on the photons. Therefore, it was not possible to use the $\eta \rightarrow \gamma\gamma$ decay to study the $\gamma n \rightarrow n\eta$ reaction, but the $\eta \rightarrow 3\pi^0 \rightarrow 6\gamma$ decay was used. For the latter decay, photon hits in the FP and/or in the MT had a reasonable probability (in the meantime the CBB has been upgraded for timing information, so that this restriction does not any more apply for future measurements). In order to avoid systematic uncertainties in the comparison of reactions off quasifree neutrons and protons, the trigger generation for η production in coincidence with recoil neutrons and recoil protons was done identically, i.e. signals from the scintillating fiber detector were not used in the trigger decision.

The trigger was composed of two levels. Details are shown in Table 1. The first level used only signals from the detectors with timing information, i.e. FP and MT, which were available in less than 400 ns. For that purpose, the MT was subdivided into four logical sectors, each comprising 54 BaF₂ modules, as shown in Fig. 1 (right hand

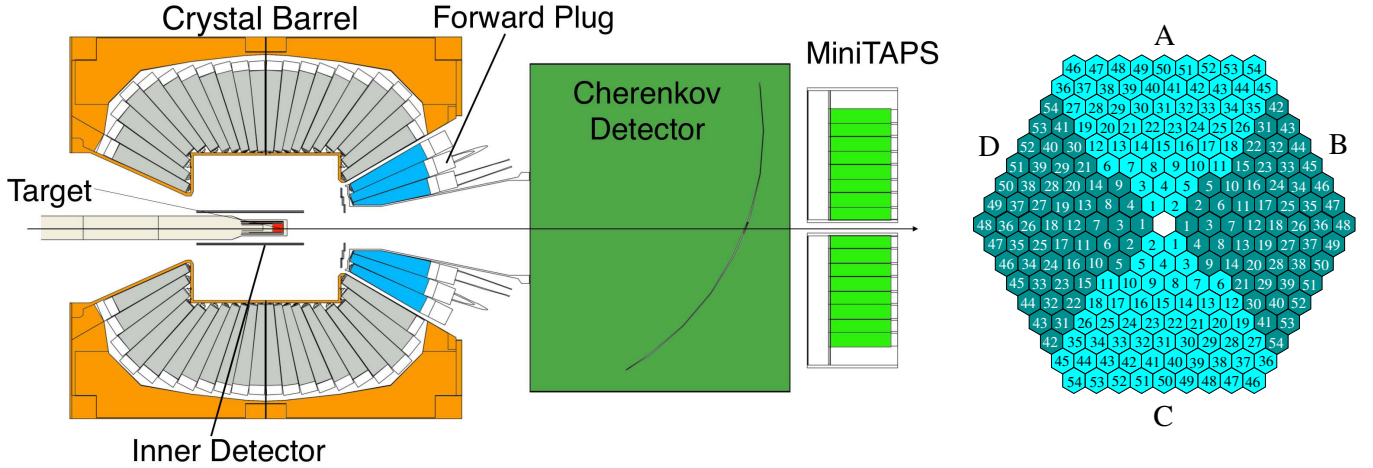


Fig. 1. Left hand side: detector setup of the CBELSA/TAPS experiment at ELSA in Bonn with the main detectors Crystal Barrel and MiniTAPS and the Cerenkov detector in between. Right hand side: front view of the MiniTAPS detector with the four logical sectors.

First Level Trigger		Second Level Trigger	
MiniTAPS	FP	FACE	
1	-	≥ 2	(≥ 3)
-	1	≥ 2	(≥ 3)
1	1	≥ 1	(≥ 2)
-	≥ 2	≥ 1	(≥ 2)
≥ 2	-	≥ 1	(≥ 2)
1	≥ 2	-	(≥ 1)
≥ 2	1	-	(≥ 1)
≥ 2	≥ 2	-	-

Table 1. Three (four) particle hardware trigger, which was used for the present work. For MiniTAPS, the numbers represent the number of hit sectors, and for the FP and FACE, the number is the cluster multiplicity. For the first level trigger, the two columns are connected with a logical ‘and’, for the second level trigger the first column is for the measurement with the unpolarized target and the second column (in brackets) for the double polarization experiment.

side). Each sector for which at least one module responded with a signal above the threshold of the discriminators (leading edge type, LED) was considered as activated. The first ring of modules closest to the beam pipe was not allowed to trigger because of the high count rates in these modules. For events with only one segment hit, the LED threshold was 100 MeV for all rings and for events with two or more segments responding it was 120 MeV for the second ring and 80 MeV for all other rings. The FP was equipped with a Cluster-Finder (CF) algorithm which combines neighboring groups of responding crystals into a cluster and delivered the number of clusters per event. Only modules with signals above a threshold of 30 MeV were considered. The first level trigger was then derived from the cluster multiplicity in the FP and the sector multiplicity in the MT as summarized in Table 1. One should note, that the LED thresholds were only relevant for the generation of the trigger, not for the readout of detector elements. When a valid trigger had been generated MT

modules were read out when the additionally available constant fraction discriminator (CFD) had a signal above threshold. These thresholds were set below 17 MeV for the two inner-most rings and below 10 MeV for the others.

The second level of the trigger included the cluster multiplicity in the CBB, which was determined with the Fast Cluster Encoder (FACE) based on cellular logic. It delivers the number of clusters in the CBB in a time range of $\approx 6 \mu\text{s}$, which is not fast enough for the first level trigger. The threshold energy for clusters was set to 20 MeV. The required CBB multiplicities in dependence of the first level trigger are given in the right hand side of Table 1. The combined multiplicity from the FP, MT, and CBB was required to be three for the measurements with the unpolarized target and four for the butanol target. A total multiplicity of four also allows to investigate events from double π^0 production, which were measured simultaneously. A cluster multiplicity of three allows to also include four-photon events where two photons hit the same sector in TAPS. For the measurement with the polarized target, a stricter trigger was used in order to increase the η count rate. During the data analysis, an additional software trigger was imposed, which required that the hardware trigger condition was already satisfied by hits assigned as photons (no trigger contribution from recoil nucleons). This was done in order to minimize systematic effects in the comparison of reactions off quasifree protons and off quasifree neutrons. Recoil protons are much more likely to generate clusters above the respective trigger thresholds than recoil neutrons and it is difficult to simulate the trigger probability of neutrons.

3 Data Analysis

The calibration procedures for the different detector components are discussed in much detail for a slightly different setup (no FP, forward region covered by a larger

TAPS wall) in [36,65]. The higher level analysis (reaction identification, extraction of cross sections, elimination of effects from Fermi smearing, etc.) was done analogously to the one reported in [39,40] for the data measured at MAMI. All analysis steps will be briefly summarized in this section.

3.1 Event Selection

A coincidence analysis for the FP and MT was performed for all events and the corresponding time spectra are shown in Fig. 2. Non-coincident events were rejected by cuts on these time spectra. However, the background level of random coincidences was very low and not visible on a linear scale. The coincidence-time resolution for two hits in the MT detector was 0.51 ns (FWHM), whereas, due to the slow rise-time of the signals from the CsI crystals, the time resolution for two hits in the FP detector was only 3.28 ns (FWHM). The tagger scintillators contributed ≈ 1 ns to the time resolution between production detector and tagger so that the total time resolution was ≈ 1.25 ns for coincidences between MT and tagger and ≈ 2.5 ns between FP and tagger (see Fig. 3).

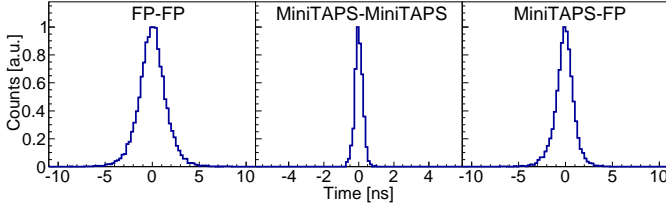


Fig. 2. Coincidence time spectra: time difference between two photon candidates in the Forward Plug (FP-FP), in MiniTAPS, and one photon candidate in MiniTAPS and one in the Forward Plug.

For each event, all electron hits in the tagger were registered during the coincidence window. The corresponding time spectra for coincidences between the MT and tagger and the FP and tagger are shown in Fig. 3. Random coincidences were first suppressed by a cut on the prompt peak and residual background below the prompt peak was removed by a sideband subtraction. The colored areas shown in Fig. 3 do not show the actual size of the windows. For this analysis, the prompt peak region (Δp) was chosen between -3 and 6 ns for MiniTAPS-Tagger and -4 and 9 ns for FP-Tagger. Events in that region had a weight of 1, whereas hits in the sidebands (between -400 and -100 ns or between 100 and 400 ns) had a weight of:

$$w = \frac{\Delta p}{\Delta R_1 + \Delta R_2}, \quad (3)$$

where ΔR_1 and ΔR_2 were the widths of the two background windows (blue). For the analysis with at least one hit in the MT, the timing from the MT was used. For other events, the timing from the FP was used. The slightly asymmetric shape of the peaks shown in Fig. 3 is

caused by massive particles misidentified as photons and disappears after the subsequent analysis steps of particle identification. The 2 ns modulation visible in the random background is due to the 500 MHz acceleration field of ELSA. The average background level in Fig. 3 is 10% for FP-tagger coincidences and 7% for MT-tagger coincidences and this is further reduced by the subsequent analysis cuts.

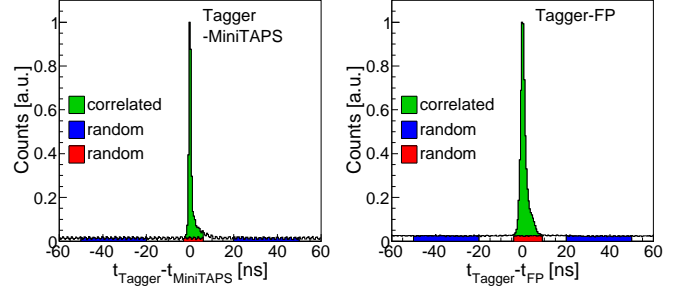


Fig. 3. Coincidence time between a hit in the tagger and the calorimeter (MiniTAPS or FP). The random hits (red) below the peak (green) were subtracted by a sideband (blue) analysis (colored areas are not actual size).

3.2 Particle Identification

In the first step of the particle identification, clusters (i.e. connected groups of activated detector modules) in the calorimeters were classified as ‘charged’ or ‘neutral’ with the help of the charge-sensitive detectors. In the CBB, clusters were identified as ‘charged’ when an intersection point of at least two fibers in the inner detector could be reconstructed, and the angular difference between the cluster in the CBB and the reconstructed hit in the inner detector was less than 12° in azimuthal and 30° in polar angle (the angular cuts were motivated by the construction and properties of the inner detector [62]). The energy threshold per fiber was set to approximately 150 keV. Clusters in the FP were regarded as ‘charged’ when at least two overlapping vetoes in the angular range of $\Delta\phi \leq 14^\circ$ and $\Delta\theta \leq 15^\circ$ were hit. No energy threshold could be set in the analysis for the FP vetoes since the energy information was not read out. In the MT, a cluster was ‘charged’ when at least one veto in front of all crystals belonging to the cluster was hit and the energy threshold was approximately 100 keV. Clusters with a single hit in the FP veto or a hit in only one of the layers of the inner detector were regarded as ‘neutral’, but rejected for the analysis of recoil neutrons.

The next step in the analysis procedure was the assignment of the events to event classes according to the number of charged (c) and neutral (n) hits. Three different event classes were defined since the $\eta \rightarrow 3\pi^0 \rightarrow 6\gamma$ decay was detected either in coincidence with the recoil proton (σ_p), the recoil neutron (σ_n), or without any condition on

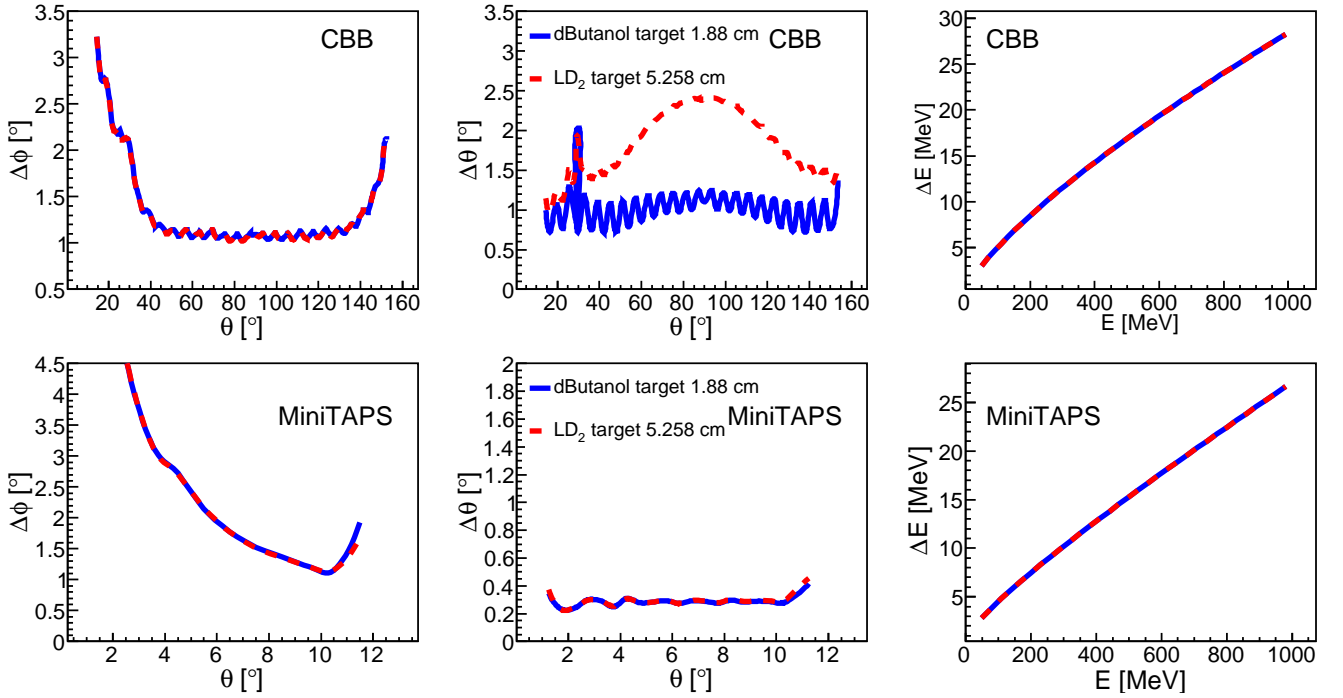


Fig. 4. Energy and angular resolutions (σ) for photons using the deuterium (red) and deuterated butanol (blue) target. The resolutions are shown for the CBB including the FP (top row) and the MT detector (bottom row).

the recoil nucleon (σ_{incl}):

$$\begin{aligned} \sigma_p: & \quad 6n \text{ and } 1c \\ \sigma_n: & \quad 7n \\ \sigma_{\text{incl}}: & \quad 6n \text{ or } 7n \text{ or } (6n \text{ and } 1c) \end{aligned} \quad (4)$$

The inclusive channel was only analyzed for the unpolarized data.

To combine the six decay photons to three π^0 , a χ^2 analysis was performed by comparing the invariant mass of two photons, $m_k(2\gamma)$, to the nominal π^0 mass, m_{π^0} :

$$\chi^2 = \sum_{k=1}^3 \frac{(m_k(2\gamma) - m_{\pi^0})^2}{(\Delta m_k(2\gamma))^2}, \quad (5)$$

where $\Delta m_k(2\gamma)$ is the uncertainty of the invariant mass, which depends on the uncertainty of the deposited energy ΔE , and the uncertainties of the azimuthal and the polar angle of the two photons, $\Delta\phi$ and $\Delta\theta$, respectively. The latter were determined using Monte Carlo simulations of isotropically distributed photons and are shown in Fig. 4 for the two different targets. It is clearly visible that the polar angle resolution for the CBB is significantly better for the short (1.88 cm) deuterated butanol target than for the longer (≈ 5.3 cm) deuterium target. The poor resolution around $\theta = 30^\circ$ was caused by the transition region between the CBB and the FP. For the MT, polar-angle resolution is identical for both targets (target extension in z -direction does not matter) and better than for the CBB (due to the larger distance). The energy resolution is similar for both calorimeters.

For events with six neutral hits or six neutral and one charged hit, the χ^2 analysis was only used to combine the

six neutral hits to the three most probable pairs corresponding to the intermediate π^0 mesons. The event was accepted when the combination with the best χ^2 has invariant masses of all three pairs of photons between 115 - 156 MeV (the mass of the π^0 is 135 MeV). The same selection was used for events with seven neutral hits and in this case, the neutral hit, which was not assigned to a π^0 decay photon, was treated as a neutron candidate. The χ^2 values were only calculated to find the best combination, no cuts were applied based on the χ^2 values.

The quality of the reconstruction is reflected in the resulting confidence level distributions, which are shown in Fig. 5. The results for the liquid deuterium target (results

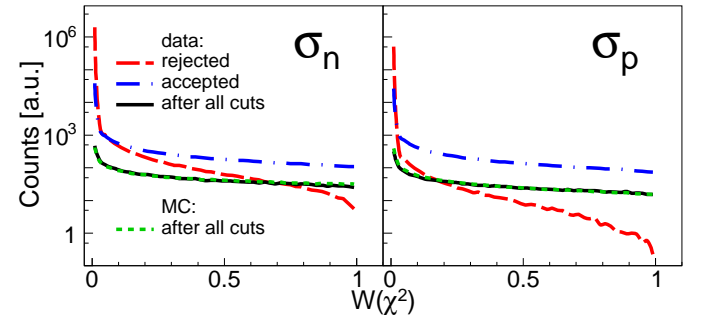


Fig. 5. Confidence level from the χ^2 test for the event classes σ_n (left) and σ_p (right). The distributions for the accepted (blue dash-dotted) and rejected (red long dashed) events are shown together with the histograms for the final events selection (black solid) for experimental data and MC events (green short dashed).

from the butanol target are very similar) are shown separately for events in coincidence with neutrons and protons. The confidence levels for the final event selection after all cuts are essentially flat and in good agreement with the Monte Carlo simulations, which indicates that the resolutions for the χ^2 test were determined realistically.

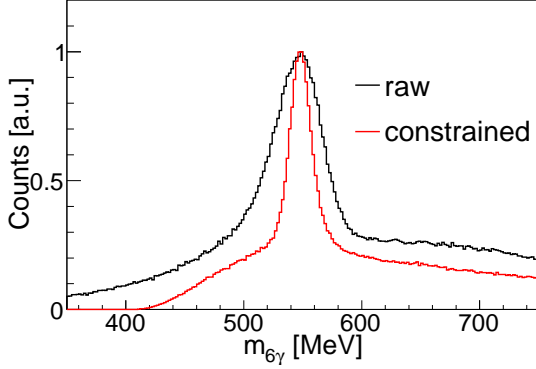


Fig. 6. Raw invariant mass $m_{\gamma\gamma}$ (black) and invariant mass after constraining (red) as shown in Eq. 6.

Subsequent to the assignment of the photon pairs to pion decays, the mass of the π^0 meson was used to improve the experimental resolution. This was done by replacing the measured energies E_{γ_1, γ_2} by E'_{γ_1, γ_2} with:

$$E'_{\gamma_1, \gamma_2} = \frac{m_{\pi^0}}{m_{\gamma_1 \gamma_2}} E_{\gamma_1, \gamma_2}, \quad (6)$$

where m_{π^0} is the nominal mass of the pion and $m_{\gamma_1 \gamma_2}$ is the invariant mass of the photon pair. Since the relative angular resolution for photons is better than the energy resolution (see Fig. 4), the angular uncertainties were neglected. This reconstruction improved the resolution for the subsequent invariant and missing mass analyses of the η mesons significantly (see Fig. 6). This improvement is in particular important for the rejection of background in the missing mass analysis.

Particle identification up to this stage used only the information from the charged particle detectors and the χ^2 analysis. Additional information was available for the MT detector from a pulse-shape analysis (PSA), time-of-flight (ToF) versus energy, and $E - \Delta E$ analyses which helped to distinguish between photons, neutrons, and protons in the angular range where most recoil nucleons were detected. These analyses are discussed in Sec. 3.4, because most of them only confirmed the correct assignment of particle types and were not used for additional cuts.

3.3 Reaction Identification

Quasifree η production was identified by subsequent cuts on the invariant mass of the six decay photons and on the reaction kinematics. For the latter, the coplanarity of the η - nucleon system and the missing mass of the reaction was investigated.

The coplanarity analysis was only possible for events for which the recoil nucleon was detected (not for the inclusive analysis). The coplanarity angle is the azimuthal angle difference between the recoil nucleon, ϕ_N , and the η meson, ϕ_η . It was calculated with the following equation:

$$\Delta\phi = \begin{cases} \phi_\eta - \phi_N, & \text{if } \phi_\eta - \phi_N \geq 0 \\ 2\pi - |\phi_\eta - \phi_N|, & \text{if } \phi_\eta - \phi_N < 0. \end{cases} \quad (7)$$

Due to four-momentum conservation, the angular difference has to be equal to 180° . Since the ϕ angle is independent of a boost in the z -direction, this is valid in the center of mass (cm) frame, as well as in the laboratory frame. However, for quasifree nucleons, the Fermi motion causes a smearing of the angle and hence, a broadening of the peak. The corresponding spectra filled right after the χ^2 selection and the PSA cut are shown in Fig. 7 for the deuterium target (colored symbols) and the deuterated butanol target (black symbols). The peaks for the data with the deuterated butanol target are broader than for the deuterium target due to the larger Fermi motion in the carbon nuclei. Since the kinematics has a strong influence on the width of the peak, the cut position of $\pm 2\sigma$ with respect to the mean position was determined for different bins of incident photon energy and $\cos(\theta_\eta^*)$ (θ_η^* is the polar angle of the η meson in the beam-target cm system assuming the initial state nucleon at rest). The same cuts were applied to the data taken with the deuterium and the deuterated butanol target.

Within the selected cuts, the experimental data from the deuterium target and the simulated spectra (solid line) are in agreement. The deviations at higher energies are mainly caused by $\eta\pi$ background. Problematic are in particular the $n\pi^+\eta$ and $p\pi^-\eta$ final states when the single charged pion escapes detection. Most of this background was later removed by the missing mass cut, which was not applied to these spectra.

For further identification of the reaction, the mass M of the recoil nucleon (treated as missing particle also when detected) was calculated from:

$$M = \sqrt{(E_\gamma + m_N - E_\eta)^2 - (\mathbf{p}_\gamma - \mathbf{p}_\eta)^2}, \quad (8)$$

where E_γ and \mathbf{p}_γ are the energy and the momentum of the incident photon beam, respectively, E_η and \mathbf{p}_η are the energy and the momentum of the η meson, respectively, and m_N is the nucleon mass. Subtracting the nominal nucleon mass from Eq. 8, i.e. $\Delta M = M - m_N$, should result in a peak at zero, which is visible in Fig. 8 for different incident photon energies. Shown are the missing mass spectra for the deuterium target (red and blue symbols), the corresponding simulated line shape (black line), and spectra for the deuterated butanol target (black circles).

In Eq. 8, it is assumed that the initial state nucleon is at rest, thus, Fermi motion has a large influence on the missing mass spectra. The Fermi motion not only induces a broadening of the peak, but also causes a slightly asymmetric shape of the peak at low incident photon energies. This is caused by the fact that close to threshold,

Fermi momenta antiparallel to the photon beam are favored since they correspond to larger cm energies.

Above $E_\gamma \simeq 800$ MeV, the background on the right hand side of the peak was mainly caused by $\gamma p \rightarrow \eta \pi^+ n$ and $\gamma n \rightarrow \eta \pi^- p$ reactions, when the charged pion was not registered. Also, the $\eta \pi^0$ final state contributed to the

background when both π^0 decay photons were not identified, which, however, is rare. This background was rejected by applying an energy and $\cos(\theta_\eta^*)$ dependent missing mass cut of $\pm 1.5\sigma$ with respect to the mean position of the peak (indicated by the vertical lines in the figure). For incident photon energies below 800 MeV, the events were accepted in the range between -100 MeV and $+100$ MeV. Here, one should note that in this missing-mass range no significant background is visible for the liquid deuterium target, the simulated line shapes agree with the experimental data. Additional background in this range for the butanol target is due to reactions with nucleons bound in the carbon (oxygen) nuclei, but this is later on subtracted (see Fig. 22).

Finally, the invariant mass of the η meson was reconstructed from the six decay photons:

$$m_{6\gamma} = \sqrt{\left(\sum_{i=1}^6 E_{\gamma_i}\right)^2 - \left(\sum_{i=1}^6 \mathbf{p}_{\gamma_i}\right)^2}, \quad (9)$$

where E_{γ_i} is the energy of the i -th photon and \mathbf{p}_{γ_i} is the corresponding momentum.

Typical spectra are shown in Fig. 9 for the different event classes. All measured data are in good agreement

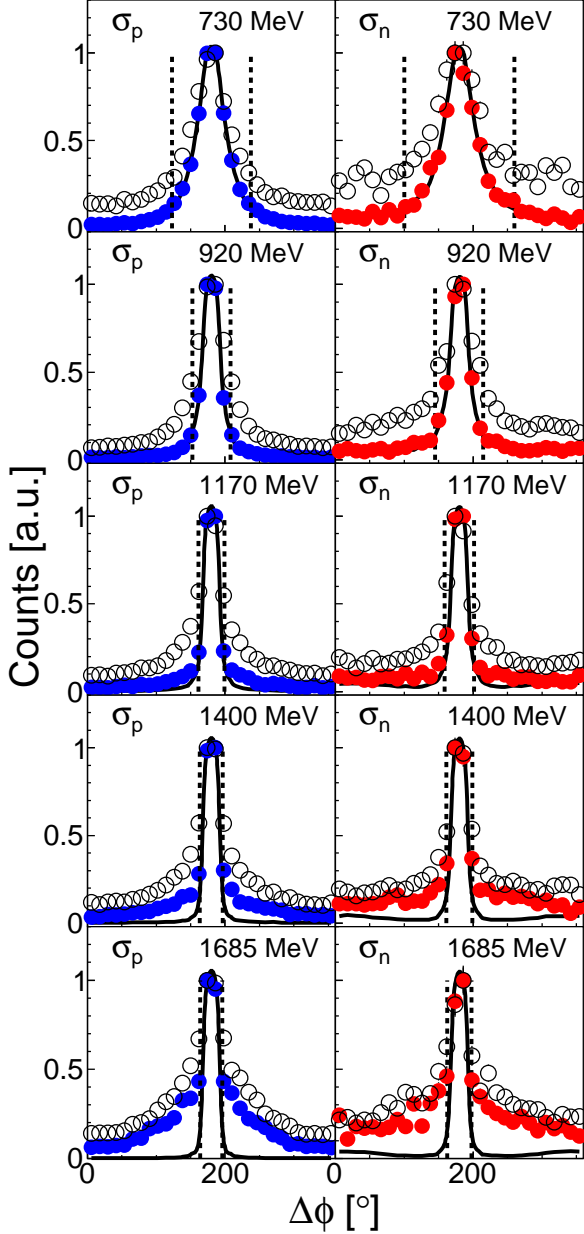


Fig. 7. Azimuthal angle difference between the η meson and the nucleon for the σ_p and σ_n event classes for five different bins of incident photon energy (energy indicated in the figure). Shown are the distributions for the deuterium target (colored solid symbols) and for the deuterated butanol target (black open) symbols. The line shape from the simulation with the deuterium target is shown as a black solid line and the cut positions ($\pm 2\sigma$) are indicated by the dashed vertical lines.

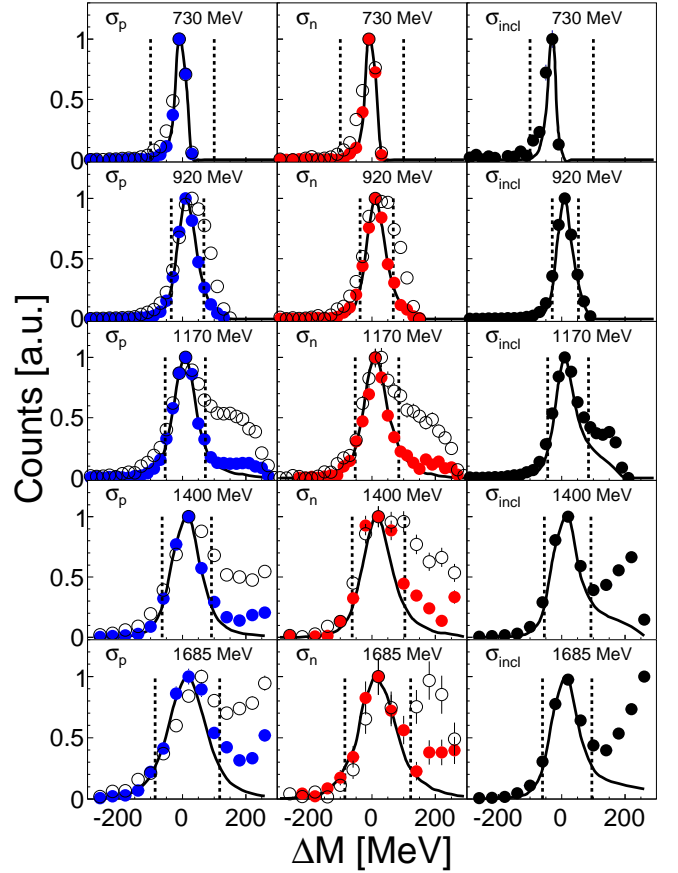


Fig. 8. Missing mass of the recoil nucleon ΔM for the event classes σ_p , σ_n and σ_{incl} for different incident photon energies (indicated in the figure). Same labeling as in Fig. 7.

with the simulated line shapes and almost free of background. Since the width and position of the invariant mass peak is essentially independent of the incident photon energy, the spectra were integrated in the invariant mass range between 500 and 600 MeV, as indicated by the vertical lines in the figure.

The spectra shown in Figs. 7, 8, and 9 are integrated over the whole angular range, however, the actual analysis was done separately for each energy and $\cos(\theta_\eta^*)$ bin.

3.4 Additional Checks

Additional detector information was used to check the particle and reaction identification, however, only very soft or no cuts were applied to these spectra.

For the forward angular range covered by the MT wall, where a large fraction of the recoil nucleons was detected, further methods of particle identification were exploited. A clean separation of neutrons from photons was achieved in the MT by a Pulse Shape Analysis (PSA). Crystals of BaF₂ have two different scintillation light components [63, 64], one with a slow decay time ($\tau = 650$ ns) and another with a fast decay time ($\tau = 0.9$ ns). The relative intensity of these two light components is different for photons and

nucleons due to their different interaction mechanisms in matter (the fast component is more intense for photons). In the experiment, the signals from the BaF₂ crystals were integrated over a short (40 ns) and a long (2 μ s) gate. For the analysis, it is convenient to define a PSA radius r_{PSA} and a PSA angle ϕ_{PSA} using:

$$r_{\text{PSA}} = \sqrt{E_l^2 + E_s^2} \quad (10)$$

$$\phi_{\text{PSA}} = \arctan \frac{E_s}{E_l}. \quad (11)$$

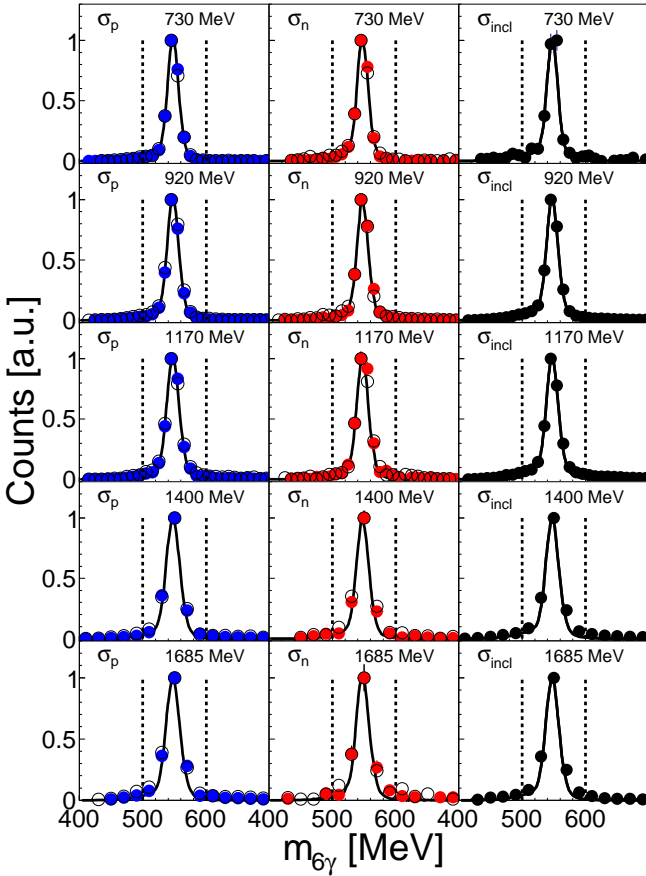


Fig. 9. Invariant mass for the event classes σ_p , σ_n , and σ_{incl} for different incident photon energies (indicated in the figure). Same labeling as in Fig. 7.

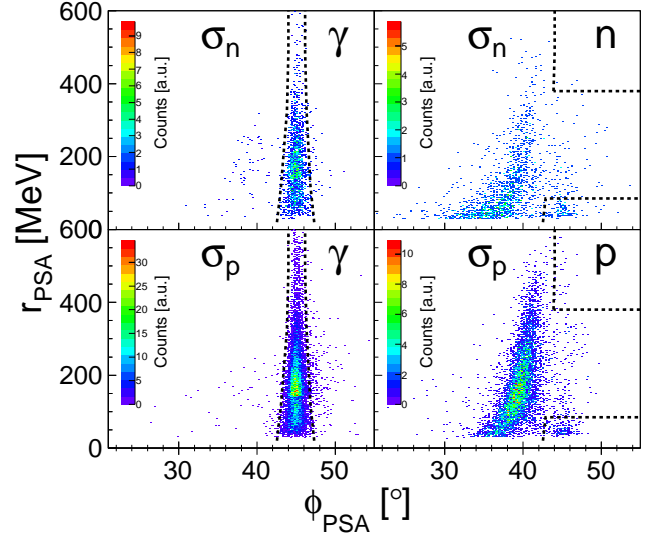


Fig. 10. Pulse shape analysis plots for the event classes σ_n (top row) and σ_p (bottom row). Shown are the PSA spectra for photon candidates (left hand side) and nucleon candidates (right hand side) after all analysis cuts were applied. The cut positions are indicated by the dashed lines. Photons are required to lie in a band of $\pm 3\sigma$ with respect to the mean position of the photon band. Nucleon candidates with $R_{\text{PSA}} < 85$ MeV or $R_{\text{PSA}} > 380$ MeV are cut away when lying closer than the 3σ line to the photon band.

Plotting the PSA angle versus the PSA radius results in the spectra shown in Fig. 10. The signals were calibrated such that the short-gate (E_s) and long-gate (E_l) responses were the same for photons, i.e. in a representation short-versus-long-gate signal photons appear on the 45° line. In contrast to the photons, the nucleons deposit less of their energy in the short gate and thus, are located at lower angles.

Having applied all other analysis cuts (i.e. χ^2 analysis, coplanarity, missing mass, and invariant mass cut) the spectra were almost free of background and soft cuts (dashed lines) were chosen to eliminate the small residual background. Photons were required to lie in a band of $\pm 3\sigma$ with respect to the mean position of the photon band. Nucleon candidates with a PSA radius smaller than 85 MeV or larger than 380 MeV were cut away when lying closer than the 3σ line to the photon band. The background in

the nucleon spectra at low PSA radii was mainly caused by electrons that did not activate the veto detectors. The cut was not applied for PSA radii between 85 MeV and 380 MeV because in this region, the nucleon band (punch-through nucleons not stopped in the MT) slightly overlaps with the photon band and real nucleons could get lost.

Since the MT provided a good time resolution, the time of flight (TOF) versus energy spectra allowed a redundant identification of photons (see Fig. 11, left hand side) which appear at a band of approximately 3.3 ns/m independent of energy. These spectra also distinguish between recoil protons and neutrons. Most protons deposit their total energy in the MT and thus appear in a band corresponding to the kinematic TOF versus energy relation for recoil nucleons, while neutrons deposit a random amount of energy in the MT. The important result is that in the TOF-versus-energy spectra for neutron candidates, no residual structure of the proton band appears, which would indicate misidentification of protons as neutrons due to inefficient charge particle counters.

Protons could be further identified by plotting the energy ΔE deposited in the MT vetoes versus the energy E deposited in the corresponding crystals. The ΔE versus E plot is shown in Fig. 12 for proton candidates. The energy resolution for these spectra was not good because the veto detectors were read out by long wave-length shifting fibers, which resulted in low photon statistics. However, they demonstrate that the background structures from electrons and charged pions (minimum ionizing peak at $E = 200$ MeV and $\Delta E = 1.5$ MeV in Fig. 12) were eliminated by the other analysis cuts.

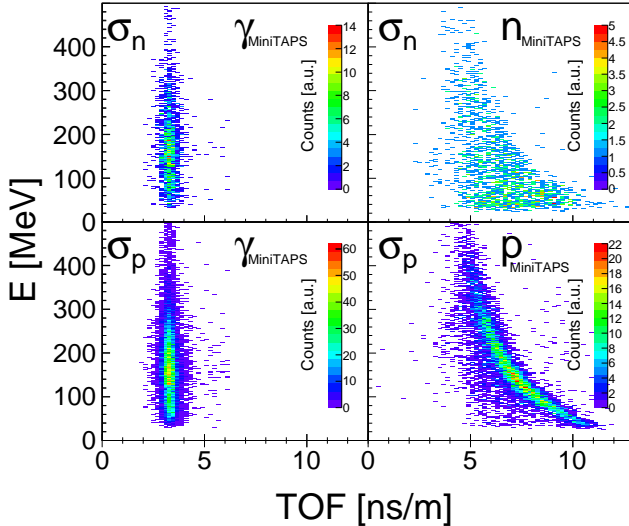


Fig. 11. TOF versus deposited energy for photons (first column), neutrons (top right), and protons (left bottom) in MiniTAPS for the event classes σ_n (top row) and σ_p (bottom row).

In the CBB (including the FP), neutrons and photons could not be directly separated. However, the identification by the χ^2 analysis and further steps of reaction identification was cross-checked by the comparison of the clus-

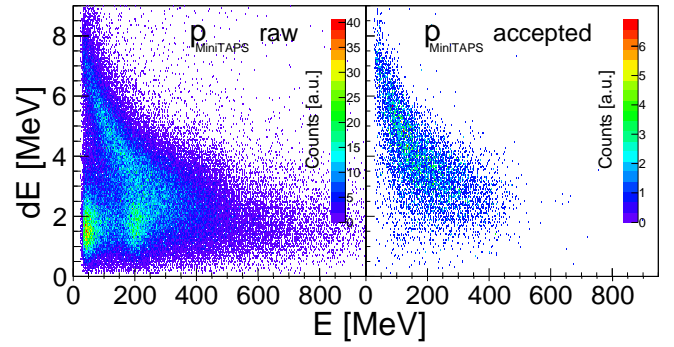


Fig. 12. ΔE versus E for proton candidates in MiniTAPS right after the event selection (left hand side) and after all subsequent analysis steps (right hand side).

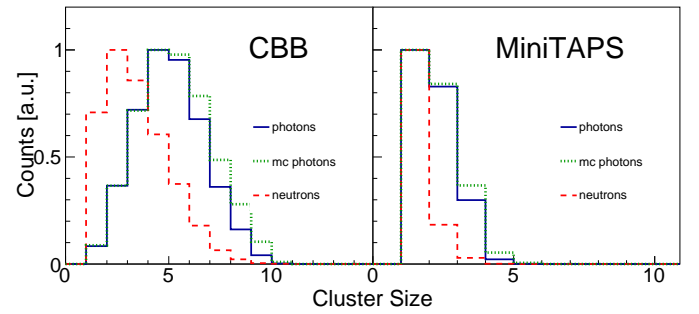


Fig. 13. Cluster multiplicity of photon and neutron candidates in the CBB (including FP) and MT. The distribution for photons and neutrons clearly differs in the CBB. For the photons, distributions from experimental data and simulation (mc) are compared.

ter multiplicity, i.e. the number of crystals per cluster of neutron and photon candidates. The corresponding distributions are shown in Fig. 13 for neutron (red) and photon (blue) candidates in the CBB and the MT. As expected, the distribution for neutrons in the CBB peaks at lower values than for photons, indicating a good separation of photons from neutrons for hits with cluster multiplicity one. However, for all other cluster multiplicities, the distributions are not well enough separated for an analysis on an event-by-event basis. The two distributions for the MT detector are more similar; however, in this case, photons and neutrons can be additionally distinguished by PSA.

3.5 Kinematic Reconstruction of the Final State

Cross sections were extracted as a function of the incident photon energy E_γ and bins of $\cos(\theta_\eta^*)$ in the beam-target cm system assuming the initial state nucleon at rest. For the two exclusive channels, where the recoil nucleon was detected, cross sections were also determined as a function of the final state invariant mass W in the η -nucleon cm system. The main advantage when using the final state invariant mass is that structures are not broadened due to Fermi motion. However, the measurement as a function of W requires the full kinematic reconstruction of the reaction, i.e. the extraction of the four-momenta of the final

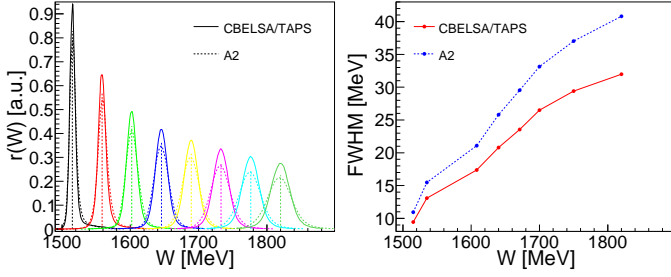


Fig. 14. Left hand side: response of the kinematic reconstruction of the final state obtained from Monte Carlo simulations with a fixed W of 1515, 1535, 1608, 1640, 1671, 1700, and 1750 MeV. Right hand side: W resolution (FWHM) for the CBELSA/TAPS setup of the current work (red dots) compared to the resolution of the A2 setup (blue dots), which was used for the work in [39, 40].

state nucleon and the η meson. As discussed in [36, 39, 40], this reconstruction is possible when the four momentum of the η -meson and the polar and azimuthal angles (direction of momentum) of the recoil nucleon are measured. The kinetic energy of the recoil nucleon is then determined by energy and momentum conservation. For nucleons in the MT, the kinetic energy can also be determined using a TOF reconstruction, however, this results in a poor W resolution [39, 40].

The W resolution achieved with the kinematic reconstruction of the final state is shown in Fig. 14. It was determined by a Monte Carlo simulation of η photoproduction reactions with a fixed invariant mass W of 1515, 1535, 1608, 1640, 1671, 1700, and 1750 MeV. The resolution obtained with the CBELSA/TAPS setup is somewhat better than the resolution of the A2 setup [40] at the Mainz MAMI accelerator, which is shown for comparison in Fig. 14. This is mainly due to a better angular resolution for the recoil nucleon because the CBB has higher granularity than the Crystal Ball in Mainz. In the region of the narrow structure in η photoproduction of the neutron, the resolution at ELSA is approximately 24 MeV FWHM, compared to the 30 MeV FWHM for the A2 setup.

The kinematic reconstruction of the final state was cross-checked by a comparison of the final-state momenta of the spectator nucleons, which also follow from the reconstruction, to the expectation from the deuteron wave function. In the approximation of a spectator - participant model without final state interactions, the final state momentum of the spectator nucleon should be equal to its initial Fermi momentum. The measured spectator momentum distributions for participant protons and neutrons are compared in Fig. 15 to the Fermi momentum distribution of nucleons bound in the deuteron extracted from the deuteron wave function of the Paris nucleon-nucleon potential [66]. For large Fermi momenta, η photoproduction can be kinematically forbidden due to overall energy conservation. Results from a simulation based on the deuteron wave function taking into account the kinematic constraints for η production are shown as dashed lines in Fig. 15 and agree quite well with the measured data.

The agreement is better for participant neutrons than for protons because for the latter energy dependent detection efficiency effects are larger and not considered in the simulation (in contrast to neutrons low energy protons do not reach the detector).

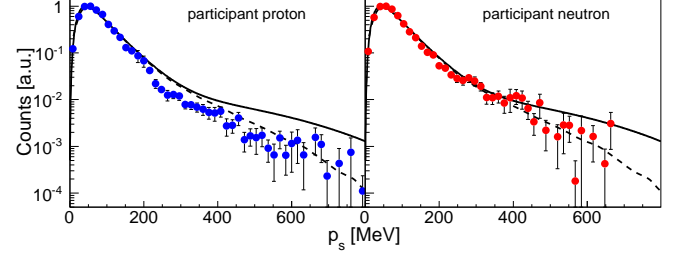


Fig. 15. Spectator momentum distribution for the participant proton (left hand side) and neutron (right hand side) in deuterium. The experimental distributions are compared to calculations using the Paris N-N potential (solid line) [66] and the corresponding kinematically allowed distribution for η photoproduction (dashed line). The y -axis is shown logarithmically.

3.6 Absolute Normalization of the Cross Sections

The cross sections were computed from the yields obtained by the integration of the final invariant mass distributions after application of all other analysis cuts. The yields were normalized to the target surface density, the $\eta \rightarrow 3\pi^0$ decay ratio $\Gamma_i/\Gamma = 32.68\%$ [47], the incident photon flux, and the detection efficiency. The different steps were done separately for each energy and $\cos(\theta_\eta^*)$ bin and are summarized in this section.

The target surface densities were 0.266 deuterons per barn for the liquid deuterium target and 0.8727 deuterons per barn (corresponding to 0.08727 butanol molecules per barn) for the deuterated butanol target. The incident photon flux was calculated from the number of electrons in the photon tagger and the tagging efficiency. The tagging efficiency is the fraction of the photons which pass through the collimator and impinge on the target. The number of photons on the target was constantly monitored by a dedicated detector system, consisting of the Gamma Intensity Monitor (GIM) and the Flux Monitor (FluMo). The rate in the GIM was 5.1 – 6.9 MHz and its efficiency was 40 – 80% (energy dependent). The efficiency was mainly reduced by dead time effects and discriminator thresholds and was corrected with the FluMo, which was calibrated at low intensities where the latter had an efficiency of 100%.

A typical spectrum for the resulting photon flux as a function of incident photon energy is shown in Fig. 16. The photon flux as a function of the final state energy W was calculated by folding the incident photon energy distribution with the Fermi momentum distribution of nucleons bound in the deuteron taken from the deuteron wave function parameterized by the Paris N-N potential [66].

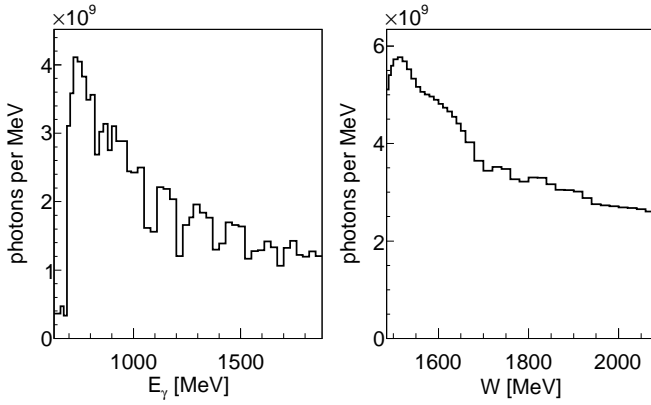


Fig. 16. Typical photon flux on the target as a function of the incident photon energy (left hand side) and the final state invariant mass (right hand side).

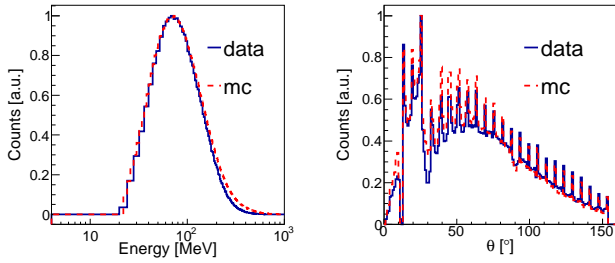


Fig. 17. Energy (left hand side) and polar angle (right hand side) distribution of the η decay photons. Compared are the spectra from experimental data and simulation (mc).

The detection efficiency was determined from Monte Carlo simulations with the GEANT3 package [67]. Events were generated with the Pluto event generator [68], including the effects of Fermi motion. They were weighted with a $1/E_\gamma$ bremsstrahlung distribution for the photon flux and cross sections for quasifree η photoproduction from the deuteron taken from [36,40]. The efficiencies were determined for each event class separately and for every energy and $\cos(\theta_\eta^*)$ bin. Typical results for nine bins of final state invariant W are shown in Fig. 18. Shown are the $\cos(\theta_\eta^*)$ distributions for the event classes σ_p (blue) and σ_n (red) for the deuterium target. The efficiencies are between 2.5% and 14% for the reaction off the proton and 1% to 2% for the reaction off the neutron.

The simulations are very reliable and precise for the detection of the η -decay photons in the detector. As an example Fig. 17 compares the measured and simulated energy and polar-angle distributions of data and simulation for the $\eta \rightarrow 6\gamma$ final state. Simulated cluster multiplicities (i.e. the number of crystals responding in one electromagnetic shower) are compared to experimental data in Fig. 13. The simulations include also the correct implementation of the experimental trigger, which was only based on photon hits as discussed in Sec. 2. In the analysis, all trigger thresholds (for measured data and simulated responses) were set by software safely above the used hardware thresholds so that no bias could arise.

The precision of the simulation is worse for recoil nucleons, in particular for charged protons detected close to the transition regions from the CBB to FP and from the FP to MT, where the results are very sensitive to small inaccuracies in the modeling of insensitive support materials and detector geometry. Therefore the detection efficiency for recoil protons was independently studied using η photoproduction off the free proton measured with a liquid hydrogen target. The number of events with detection of the recoil proton was compared to the total number of events with an identified η meson. The ratio of these two event rates corresponds directly to the proton detection efficiency. This was done for the measured data and also for the results of the Monte Carlo simulation of the hydrogen experiment. The comparison of the detection efficiency derived from the measurement and the simulated one, both as a function of proton laboratory angle θ_p and proton kinematic laboratory energy T_p , was then used to extract correction factors for the simulation of the deuterium experiment. Fig. 19 shows the relative proton efficiency correction as a function of θ_p and T_p . The largest deviations from unity occur in the region of the FP detector ($11.8^\circ \leq \theta_p \leq 27.5^\circ$) and the transition from the CBB to the FP around $\theta_p \approx 30^\circ$. They are most probably caused by geometrical discrepancies between simulation and experiment and the fact that no precise energy thresholds could be determined for the FP vetoes and the inner detector, since they did not provide energy information. It is shown at the right hand side of Fig. 19 that the correction significantly improves the agreement of the present results with previous measurements [22] of the free proton cross section in the energy range of the $N(1535)$ resonance. This effects are negligible at larger W .

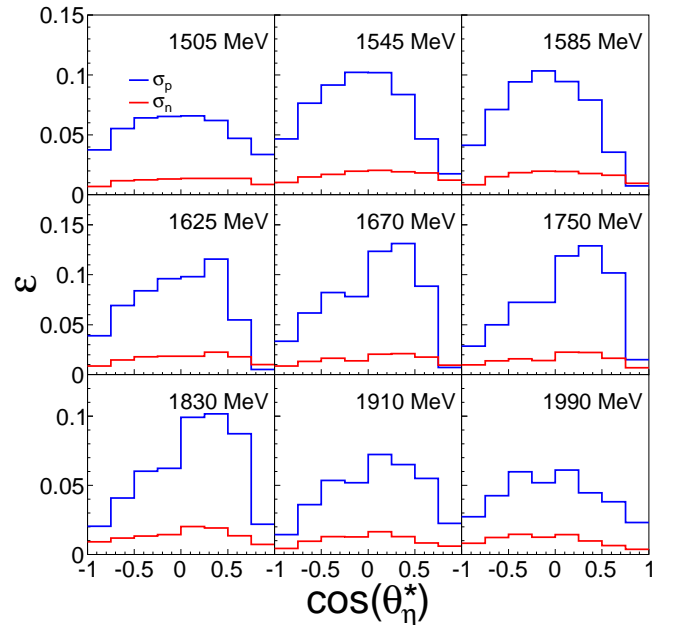


Fig. 18. Angular detection efficiencies for different bins of the final state invariant mass W (energy indicated in the figure) for the reaction off the proton (blue) and off the neutron (red).

A correction of the neutron efficiency (as in [39,40]) was not possible, since the analysis would require the unambiguous identification of the charged pion in the reaction $\gamma p \rightarrow \pi^+ n$ or $\gamma p \rightarrow \pi^0 \pi^+ n$. This was only possible in the MT detector via a $\Delta E - E$ analysis, but not in the CBB detector or in the FP detector due to the missing energy information of the charge sensitive detectors. However, analyses of data from a similar detector setup at the MAMI accelerator [39,40], for which identification of charged pions was possible, have shown that the effects for neutron detection are smaller because low energy neutrons are much less affected by passing through inactive structural material of the detector.

The determination of the unpolarized cross sections also included the subtraction of the contribution from the target cell ($2 \times 125 \mu\text{m}$ capton foil). Since no empty target measurement was available, this contribution was estimated from a measurement with a carbon target. The yields were normalized to the number of nuclei in the target windows and subtracted to obtain the final cross section values. The empty target contribution was up to 5% for the exclusive cross sections σ_p , σ_n and up to 10% for the quasifree inclusive reaction.

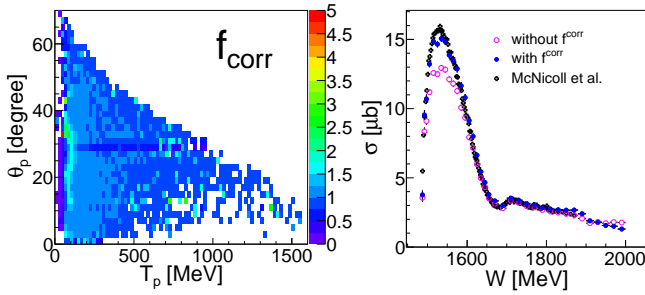


Fig. 19. Left hand side: proton efficiency correction factor f_{corr} as a function of the laboratory angle of the proton θ_p and its kinetic energy T_p . Right hand side: effect of f_{corr} on the total cross section. The cross sections from the current experiment with (blue dots) and without correction factor (magenta open circles) are compared to results on the free proton target from [22] (black stars).

3.7 Extraction of the Double Polarization Observable E and the Helicity Dependent Cross Sections

The analysis of the polarized data was done analogously to unpolarized data. The yields with helicity 1/2 and 3/2 were normalized with nucleon numbers, branching ratio, photon flux, and detection efficiency.

The statistical quality of the data measured with the polarized target was limited and not sufficient for an exclusive analysis of η production in coincidence with recoil neutrons. However, due to the low neutron detection efficiency, statistics for the measurement of the $\gamma n \rightarrow n\eta$ reaction is increased by roughly a factor of three when requiring that no recoil proton has been detected instead of requiring the detection of a coincident neutron. Since

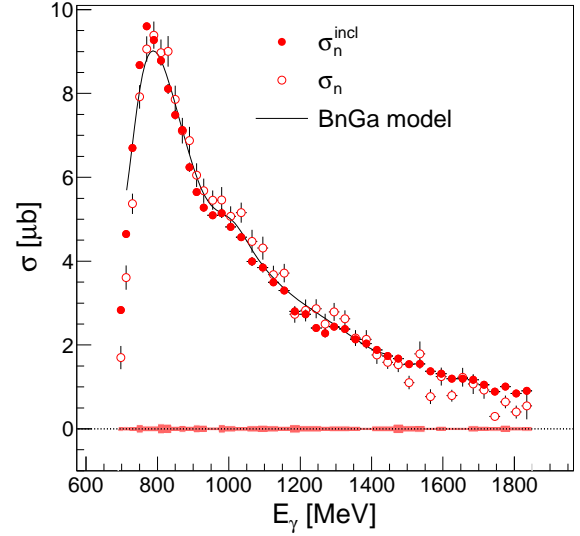


Fig. 20. Comparison of the exclusive (σ_n , open symbols) and semi-inclusive (σ_n^{incl} , closed symbols) analysis of the reaction $\gamma n \rightarrow n\eta$. Curve: Fermi folded model calculation by BnGa [51]. The systematic uncertainty of the inclusive cross section is given by the shaded area on the bottom, except 10% total normalization uncertainty.

the proton detection efficiency is typically above 95%, this type of analysis does not introduce substantial background. The analysis accepted only events for which none of the charge sensitive detectors responded. This condition enhances proton suppression because in contrast to the analysis of the exclusive proton data, hits where the proton was stopped in the charged particle detectors were also assigned to protons. The method was tested with the unpolarized data for which also the exclusive $\gamma n \rightarrow n\eta$ reaction with detection of the recoil neutron could be studied. In Fig. 20, total cross sections of the exclusive (σ_n , coincident detection of recoil neutron) and semi-inclusive (σ_n^{incl} , vetoing of events with charged hits) analyses for the $\gamma n \rightarrow n\eta$ reaction are compared. They are in quite good agreement, but the statistical quality of the inclusive data is much better. The drawback of this analysis method is that without detection of the recoil neutron, the final state kinematics cannot be reconstructed so that the effects from nuclear Fermi motion are not removed and all structures in the total cross section are significantly broadened by it.

The double polarization observable E was deduced in two different ways, corresponding to the two parts of Eq. 1. In version (1) (first part of Eq. 1), the difference of the two helicity dependent cross sections was divided by the sum of the two. This analysis required the subtraction of the unpolarized background from carbon and oxygen nuclei in deuterated butanol. This background becomes clearly visible when one compares the missing mass spectra of the sum and the difference of the two helicity states. This is shown in Fig. 21. The spectra of the helicity difference correspond only to reactions on nucleons bound in

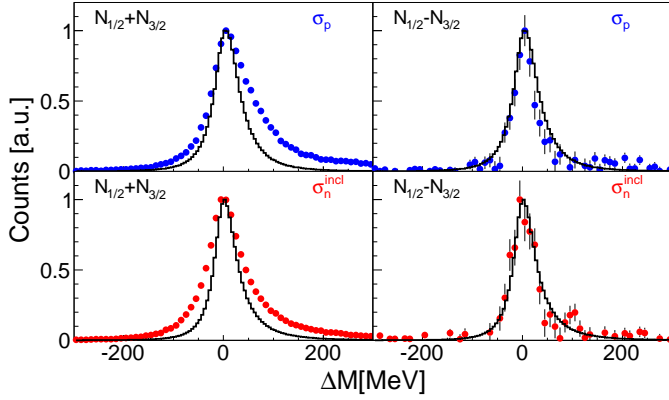


Fig. 21. Missing mass for the sum (left) and the difference (right) of two helicity states for the exclusive reaction off the proton (top) and the inclusive reaction off the neutron (bottom). The experimental spectra (dots) are compared to the simulated line shape (black line) for nucleons bound in the deuteron.

the deuteron. The background from the unpolarized nucleons bound in the heavier spin-zero nuclei drops out. These spectra agree very well with a MC simulation for the reaction on the deuteron, taking into account Fermi motion in the deuteron. The helicity sum spectra are much broader due to the unpolarized background from the heavier nuclei with larger Fermi momenta.

The contribution of the unpolarized carbon and oxygen nuclei was extracted from a dedicated measurement with a carbon foam target of a similar density as the deuterated butanol target. The data were analyzed with the same procedure as the data from the deuterated butanol and the yields were extracted and normalized with nucleon numbers, branching ratios, photon flux, and efficiencies. The normalized yields from the carbon measurement were then directly subtracted from the ones of the measurement with the deuterated butanol target. This procedure was cross-checked with a comparison of the yields from the butanol, carbon, and the liquid deuterium target, which is shown in Fig. 22. The yields from all three targets were normalized absolutely as discussed above (no free parameter). The results for the deuterium target and the carbon target were added and agree very well with the measurements of the butanol target. For the carbon subtracted spectra (not shown), the cuts indicated in Fig. 22 were applied.

In a second analysis version (version (2)), the difference of the two helicity states was normalized to the unpolarized cross section measured with the liquid deuterium target, corresponding to the second part of Eq. 1.

The helicity dependent cross sections $\sigma_{1/2}$ and $\sigma_{3/2}$ follow from the E asymmetry via:

$$\begin{aligned}\sigma_{1/2} &= \sigma_0(1 + E) \\ \sigma_{3/2} &= \sigma_0(1 - E),\end{aligned}\quad (12)$$

where σ_0 is again the unpolarized cross section measured with the liquid deuterium target. The two different extraction methods of E correspond to two solutions for $\sigma_{1/2}$ and

$\sigma_{3/2}$, which both use the unpolarized cross section σ_0 (see Eq. 12) for the absolute normalization. A third extraction method (version (3)) for $\sigma_{1/2}$ and $\sigma_{3/2}$ uses directly the two cross sections measured with the butanol target for the parallel and antiparallel spin orientations with subtraction of the carbon background and is thus completely independent of the liquid deuterium data. All versions are statistically correlated, but have a different systematic error and hence are ideally suited to crosscheck the extraction of the observables.

3.8 Systematic Uncertainties

For the unpolarized cross sections, the systematic uncertainty is composed of overall normalization uncertainties and energy and $\cos(\theta_\eta^*)$ dependent effects.

The normalization uncertainties come from the target surface density (7%, estimated from uncertainties in the measurement of the target pressure and possible changes of target length between room temperature and cooled state), the η decay branching ratio (0.8% calculated from [47]), the photon flux (8%, mainly related to the precision of the GIM and the trigger electronics), and the empty target contribution (2.5%, very conservatively estimated to be approximately half the contribution of the empty target measurement). These contributions were added in quadrature for a total uncertainty of approximately 10%.

The energy and $\cos(\theta_\eta^*)$ dependent effects result from the analysis cuts and the MC simulations, more precisely from how well the MC simulations reflect the analysis cuts. They were determined separately for each reaction channel and estimated by a comparison of the results from analyses with varying cuts on the coplanarity angle, the missing mass, and the invariant mass. Only the precisely determined incident photon energies, the parameters (angle and energy) of the η -decay photons, and the angles of the recoil nucleons enter into this cuts. In the comparison, all cuts were simultaneously varied by $\pm 3\%$.

The uncertainty of the nucleon efficiency correction was investigated similarly as in [36, 39, 40] by a comparison of the sum $\sigma_p + \sigma_n$ of the two exclusive cross sections σ_p , σ_n and the total inclusive cross section σ_{incl} measured without any conditions for recoil nucleons. The latter is independent on nucleon detection efficiencies. Since contributions from coherent production of η mesons are completely negligible [28], any deviations between $\sigma_p + \sigma_n$ and σ_{incl} result from imperfect nucleon detection efficiencies. Half of the observed differences were assigned to the uncertainty of the proton detection efficiency and the other half was assigned to that of the neutron. The uncertainties induced by the analysis cuts and the detection efficiency are shown in Fig. 23 for different final state energy bins and as a function of $\cos(\theta_\eta^*)$.

The systematic uncertainties of the double polarization observable E and of the helicity dependent cross sections $\sigma_{1/2}$ and $\sigma_{3/2}$ were determined from the uncertainty of the target polarization (5%), the photon polarization (3%), and the smoothed difference between the analysis versions (in case of three versions, the maximum difference was

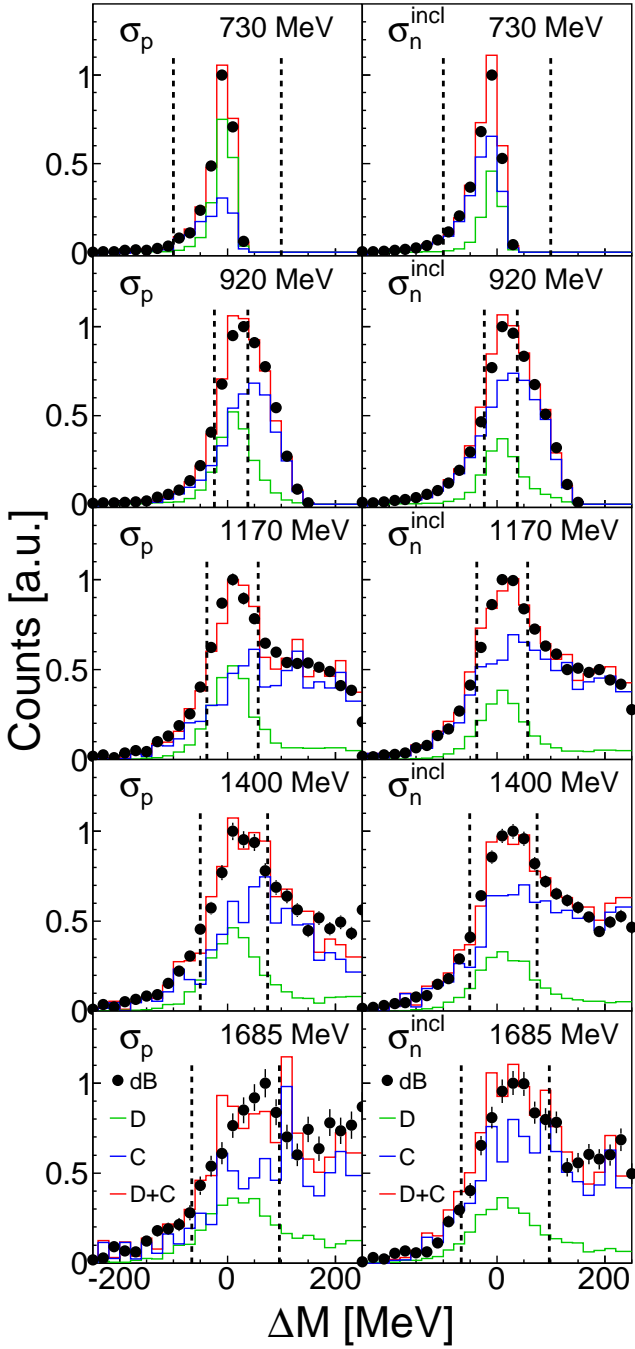


Fig. 22. Normalized missing mass spectra for the reactions $\gamma p \rightarrow p\eta$ (left hand side) and $\gamma n \rightarrow (n)\eta$ (right hand side) for five bins of incident photon energy (mean energy indicated). The data points from the measurement with the deuterated butanol target are shown in black, and the line shapes from the measurement with the carbon and deuterium target are shown in blue and green, respectively. The red line is the sum of the deuterium and the carbon data. The vertical lines are the missing mass cut positions.

taken). This estimation is reasonable since the systematics

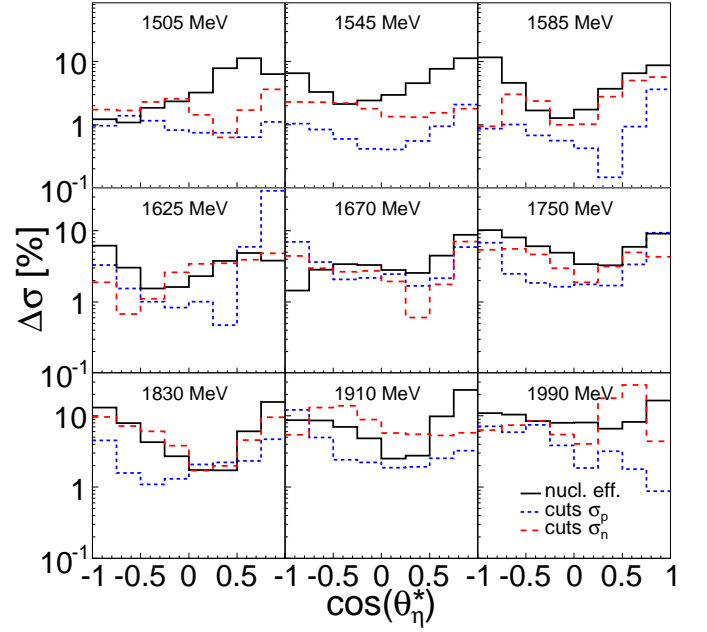


Fig. 23. Typical relative systematic uncertainties induced by the analysis cuts (proton: blue short dashed, neutron: red long dashed lines) and the nucleon efficiency (solid line, same for proton and neutron).

of all versions is independent even though the statistical errors are highly correlated.

4 Results

The results presented in this section for the unpolarized total cross section and angular distributions were obtained from the analysis of the liquid deuterium data. The analysis of the asymmetry E and the helicity dependent cross sections $\sigma_{1/2}$, $\sigma_{3/2}$ used data from the polarized deuterated butanol target, the liquid deuterium target, and the solid carbon target. The event selection and particle identification was identical for all data sets.

4.1 Unpolarized Cross Sections

The angular distributions have been extracted as discussed above in small bins of energy E (either incident photon energy E_γ or final state energy W) and for eight equidistant bins of $\cos(\theta_\eta^*)$. The total cross sections were deduced from the angular distributions via a fit of Legendre polynomials of third order:

$$\frac{d\sigma}{d\Omega}(E, \cos(\theta_\eta^*)) = \frac{q_\eta^*(E)}{k_\gamma^*(E)} \sum_{i=0}^3 A_i(E) P_i(\cos(\theta_\eta^*)), \quad (13)$$

where q_η^* and k_γ^* are the η and photon momenta in the center-of-mass frame, respectively, and $A_i(E)$ are the Legendre coefficients.

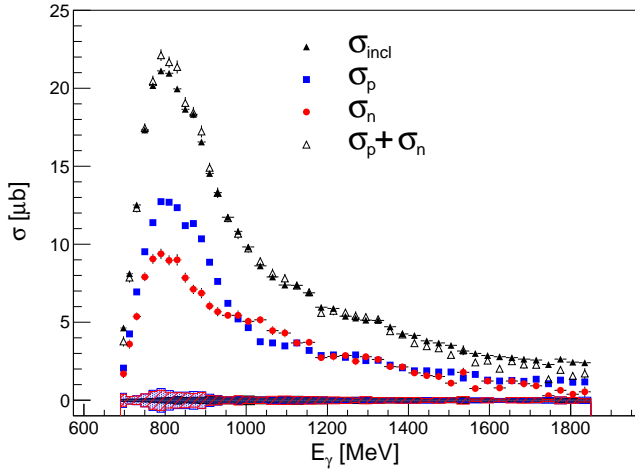


Fig. 24. Total exclusive cross sections σ_p (blue squares), σ_n (red dots), and the inclusive cross section σ_{incl} (black triangles) as a function of the incident photon energy. In addition, the sum of the two exclusive cross sections $\sigma_p + \sigma_n$ (black open triangles) is shown. The estimated systematic uncertainties are given at the bottom, except 10% total normalization uncertainty.

The total cross sections in the beam-target cm system assuming the initial state nucleon at rest are shown in Fig. 24 for the exclusive reaction on quasifree protons σ_p and on quasifree neutrons σ_n , and the sum of these two cross sections $\sigma_p + \sigma_n$ is compared to the inclusive cross section σ_{incl} . The excellent agreement between $\sigma_p + \sigma_n$ and σ_{incl} sets stringent limits for systematic uncertainties in the nucleon detection efficiencies, which would affect the proton/neutron cross section ratio.

The present results for the exclusive cross sections σ_p and σ_n as a function of E_γ are compared in Fig. 25 to previous results from CBELSA/TAPS [36] and Crystal Ball at MAMI [40]. The agreement for the proton data is excellent and for the neutron data it is good. For the latter, in the region around the narrow bump at $E_\gamma \approx 1$ GeV, the present data are somewhat in between the two previous measurements. Note that for the present data, the results from the exclusive analysis with detection of the coincident neutron are shown. The results from the analysis based on rejection of events with proton candidates (see Fig. 20) have better statistical quality (in particular above 1.4 GeV those results are certainly superior), however, for consistency we compare the data sets which have been analyzed in the same way.

The angular distributions in the beam-target cm system are summarized in Fig. 26. The angular dependence is in agreement with previous measurements [40]. Close to threshold, the distributions are almost flat and the slightly asymmetric shape is a Fermi motion effect. Around incident photon energies of 800 MeV, in the excitation maximum of the $N(1535)1/2^-$, angular distributions for the proton target show a shallow maximum and those for the neutron target show a shallow minimum around polar angles of 90° . This is a well understood effect from the in-

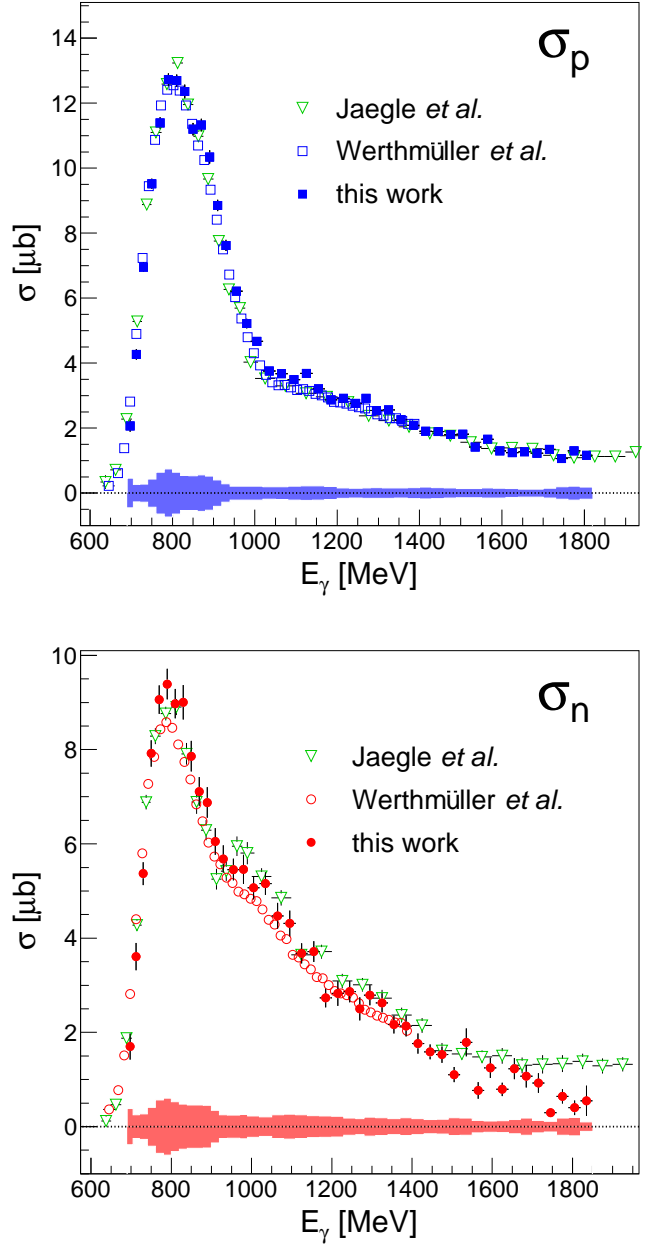


Fig. 25. Total exclusive cross sections σ_p (top) and σ_n (bottom) from this work compared to the results from [36, 40]. The estimated systematic uncertainties are given at the bottom, except 10% total normalization uncertainty.

terference between the contributions from the dominant $N(1535)1/2^-$ resonance and the tiny contributions from the $N(1520)3/2^-$ state [10, 29, 36, 40].

The total cross section as a function of the final-state invariant mass are shown in Fig. 27 for the proton (left) and the neutron (right). The systematic uncertainty (except 10% total normalization uncertainty) is indicated by the shaded histograms. The present data agree quite well with the results from Werthmüller *et al.* [40]. However, in contrast to the data as a function of incident pho-

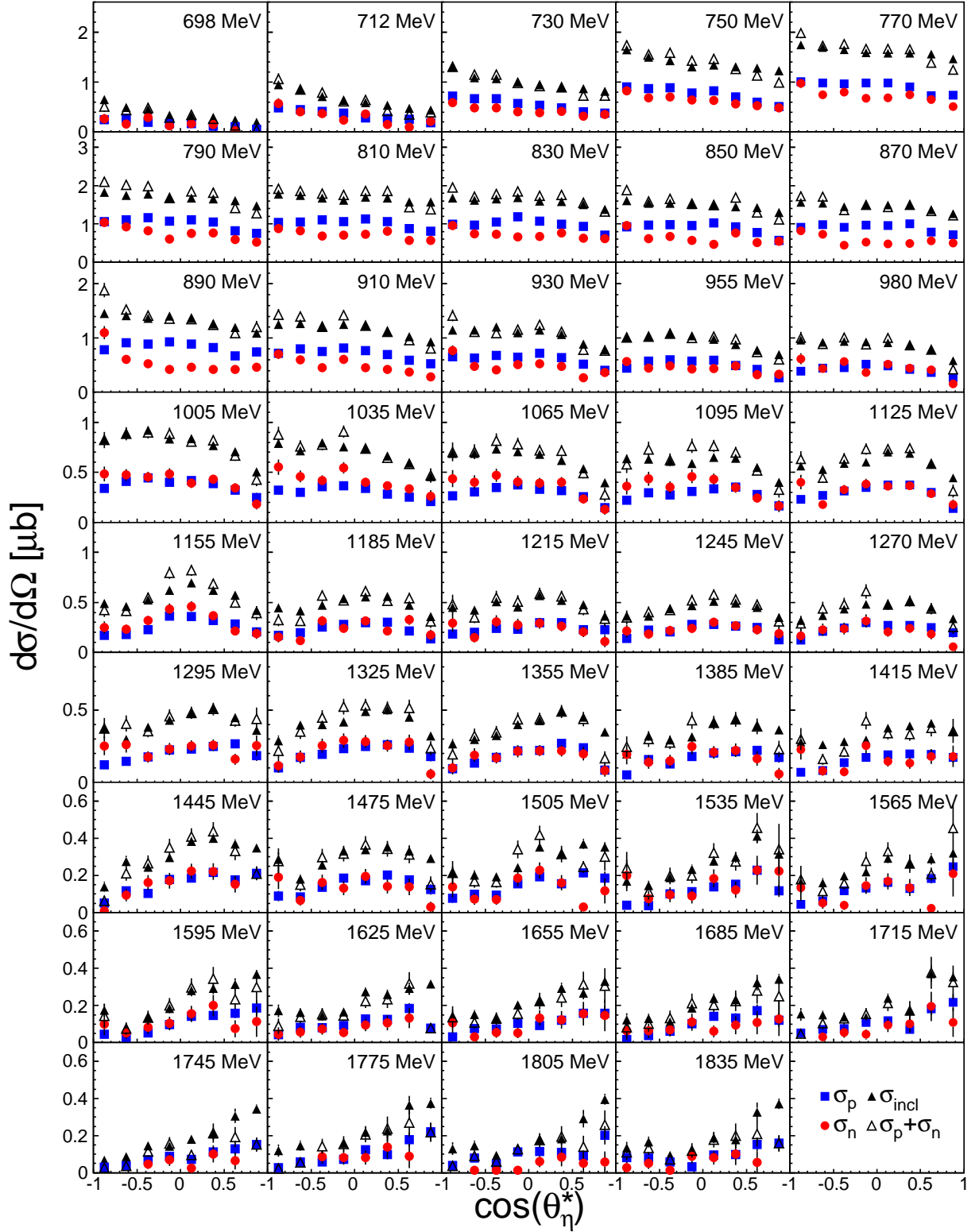


Fig. 26. Angular distributions σ_p , σ_n , and σ_{incl} for different bins of incident photon energy E_γ (mean energy indicated in the figure) as a function of $\cos(\theta_\eta^*)$ in the beam-target cm system assuming the initial state nucleon at rest. Same labeling as in Fig. 24.

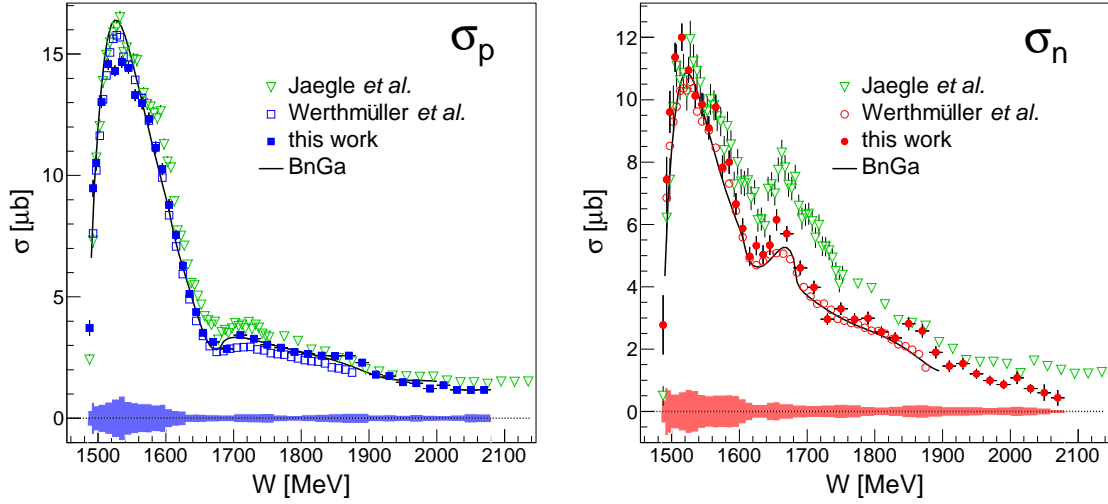


Fig. 27. Total exclusive cross sections as a function of the final-state invariant mass W for the reaction $\gamma p \rightarrow p\eta$ (left) and the reaction $\gamma n \rightarrow n\eta$ (right). Previous results from Werthmüller *et al.* (MAMI) [40] and Jaegle *et al.* (ELSA) [36], and BnGa model analyses [51] are compared. The systematic uncertainties are indicated at the bottom, except 10% total normalization uncertainty.

ton energy without kinematic reconstruction of the final state (see Fig. 25), they disagree significantly with the CBELSA/TAPS results of Jaegle *et al.* [36]. Already in [25] it was shown (by a comparison of the integrals) that the kinematically reconstructed results from [36] are inconsistent with the Fermi-motion uncorrected data from the same work, which most likely is caused by an energy dependent normalization error (the effect is similar for proton and neutron data).

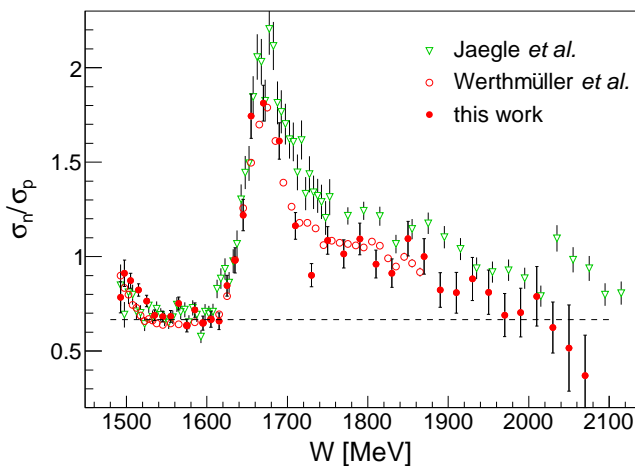


Fig. 28. Ratio of neutron and proton cross section as a function of the final state energy W . The data of this work (red dots) are compared to previous results by Werthmüller *et al.* (red circles) [40] and Jaegle *et al.* (green triangles) [36]. The dashed line marks the ratio of $\sigma_n/\sigma_p = 2/3$, which was already seen in earlier experiments [26,27,30] and is expected from the excitation of the $N(1535)1/2^-$ state.

The present results confirm this finding and the W dependent CBELSA/TAPS data from [36] are superseded by the present results. Although statistically not very significant, the structure in the neutron excitation function around $W \approx 1.67$ GeV (and also the dip-like structure in the proton excitation function at the same W), appear a bit sharper than in the previous MAMI data [40]. This is consistent with the better W resolution of the CBELSA/TAPS experiment due to the better angular resolution of the CBB compared to the Crystal Ball. Agreement with the results from the BnGa model [51] is also good because they are based on fits to the previous MAMI data [40].

The ratios of the neutron and proton cross sections from the present and previous [36,40] data are compared in Fig. 28. They are in reasonable agreement for all measurements and emphasize the rapid change above $W = 1.6$ GeV. At smaller W , the ratio is dominated by the excitation of the $N(1535)1/2^-$ state for which a $\sigma_n/\sigma_p \simeq 2/3$ scaling [26,27,30] was established by earlier experiments.

The angular distributions for different bins of final state energy W are shown and compared to previous results from [36,40] in Fig. 29 for the proton and in Fig. 30 for the neutron. Agreement between the present data and [40] is on overall quite good.

The previous data from [36] were more coarsely binned (only four data points per angular distribution) and, as discussed above, are affected by a normalization problem, but even those data show a similar angular behavior. The present data extend the W range for the angular distributions by 160 MeV. Over most of the energy range, angular distributions for proton and neutron participant nucleons are different and this is also reflected in the BnGa model results (which have been fit to the data from [40]).

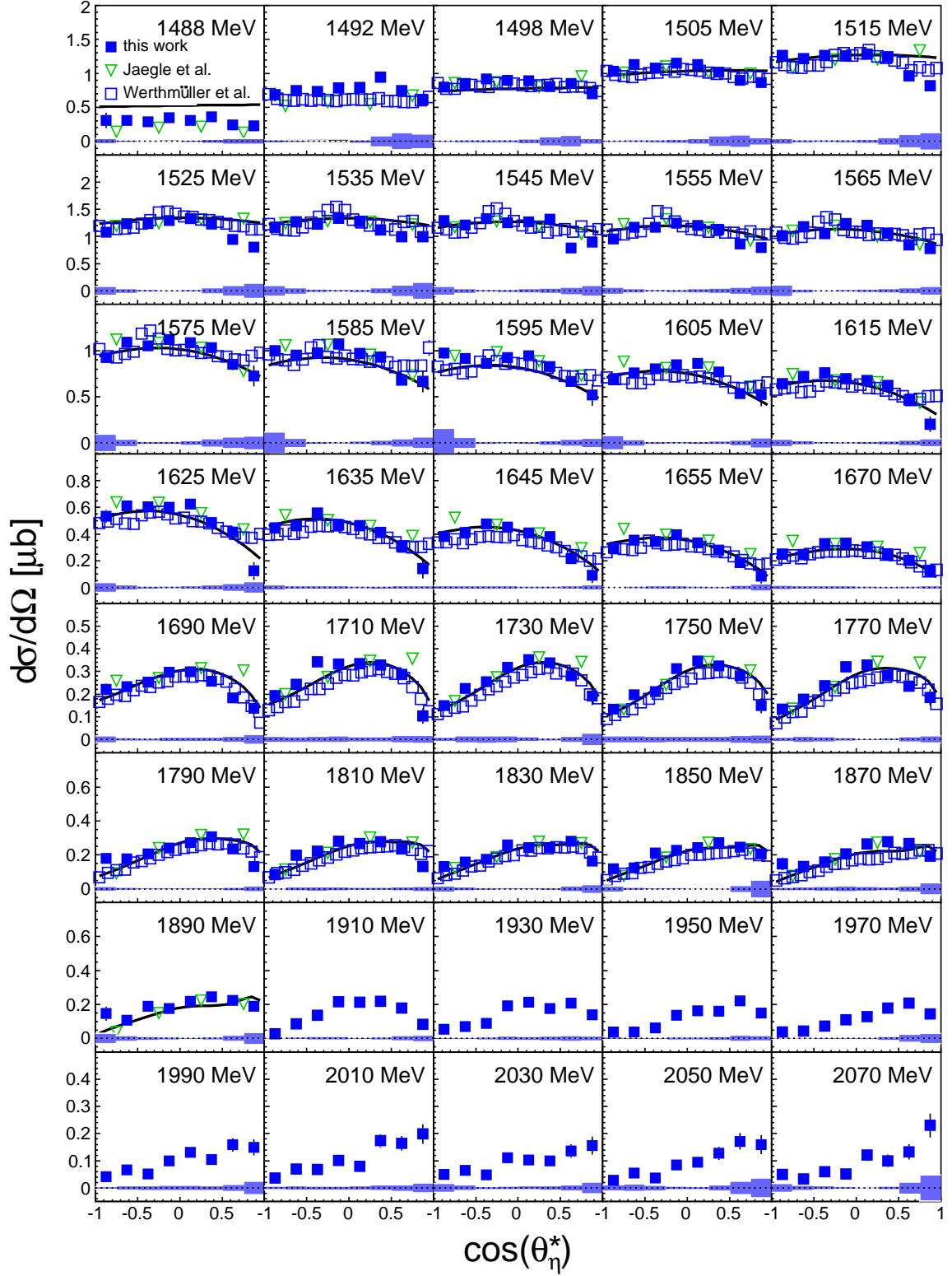


Fig. 29. Angular distributions in the η -nucleon cm system for the reaction $\gamma p \rightarrow p\eta$ for different bins of final state energy W (mean energy indicated in the figure). Blue squares: current data, open squares: results by Werthmüller *et al.* [40], green triangles: Jaegle *et al.* [36], black line: recent model analysis by BnGa [51]. The systematic uncertainties are given in bottom of each histogram, except 10% total normalization uncertainty.

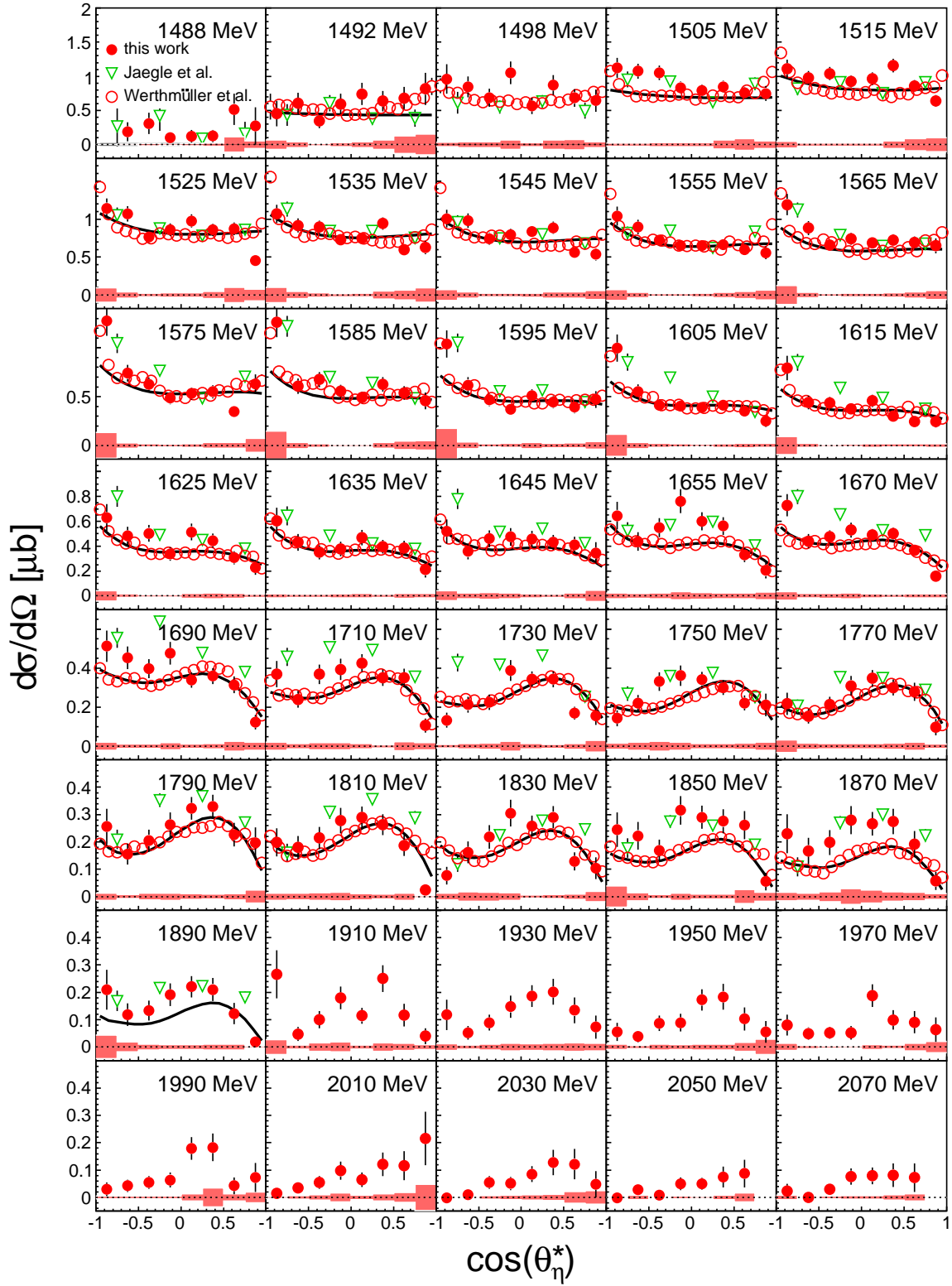


Fig. 30. Angular distributions in the η -nucleon cm system for the reaction $\gamma n \rightarrow n\eta$ for different bins of final state energy W (mean energy indicated in the figure). Red dots: current data, red circles: results by Werthmüller et al. [40], green triangles: Jaegle *et al.* [36], black line: recent model analysis by BnGa [51]. The systematic uncertainties are given in bottom of each histogram, except 10% total normalization uncertainty.

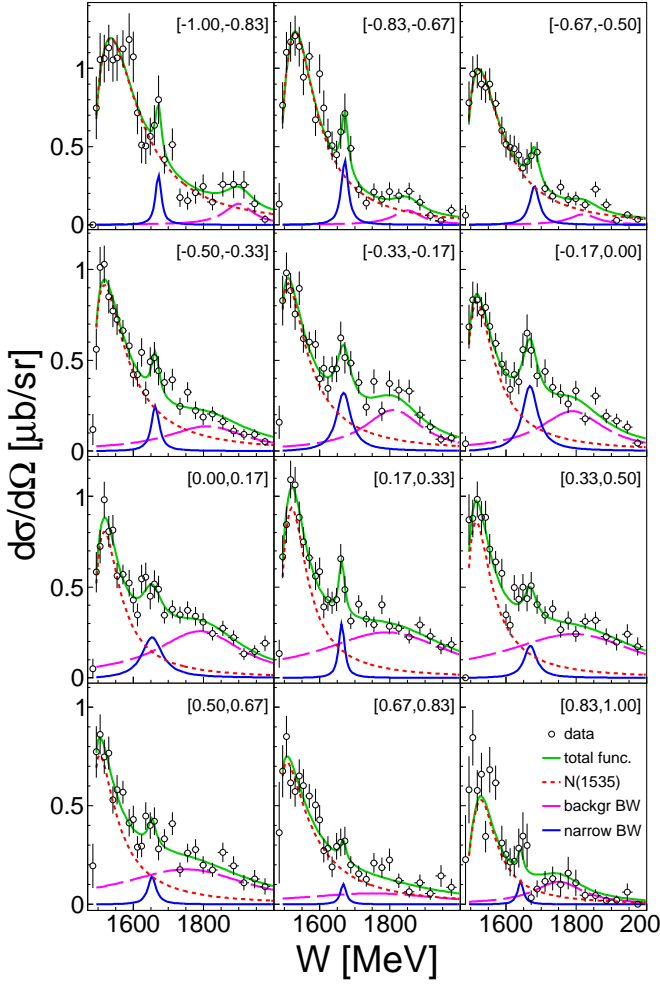


Fig. 31. Differential cross sections in the eta-nucleon cm system for the reaction $\gamma n \rightarrow n\eta$ for 12 angular bins ($\cos(\theta_\eta^*)$ range indicated in the figure). The total fit function (green line) is the sum of a energy dependent BW (red dashed line), a broad background BW (magenta long dashed), and a narrow BW (blue) function.

A fit with the same phenomenological ansatz as in [40], i.e. with three Breit-Wigner functions: one for the S_{11} wave, one for the narrow structure, and a third subsuming non-resonant backgrounds and other partial waves, is used for a comparison to previous results. The fit yields a position of 1667 ± 2 MeV, a width of 35 ± 3 MeV, and an electromagnetic coupling of $\sqrt{b_\eta} A_{1/2}^n = 13.4 \pm 2 \cdot 10^{-3} \text{GeV}^{1/2}$ for the narrow structure. Accounting for the experimental resolution, a width of only 23 ± 2 MeV is found. Table 2 summarizes all results of the fit and compares them to previous results. Note that the results for the S_{11} wave are not corrected for experimental resolution so that in particular the width Γ deviates between the present results and [38, 40].

In order to investigate the angular dependence of the narrow structure in more detail, differential cross sections were extracted for 12 angular bins. The distributions are

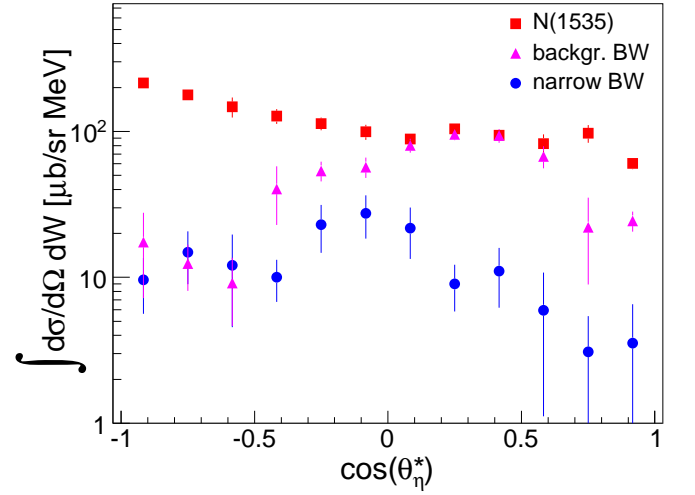


Fig. 32. Integral of the angular fits to the differential cross sections from Fig. 31. Red squares: integral of the $N(1535)$ BW, magenta triangles: integral of the background BW, and blue dots: integral of the narrow BW.

		W_R [MeV]	Γ [MeV]	$\sqrt{b_\eta} A_{1/2}^n$ [$10^{-3} \text{GeV}^{1/2}$]
narrow structure	this work	1667 ± 3	35 ± 3 (23 ± 2)	13.4 ± 2
	[38]	1670 ± 1	50 ± 2 (29 ± 3)	12.3 ± 0.8
	[36]	1663 ± 3	25 ± 12	-
$N(1535)$	this work	1525 ± 2	146 ± 13	82 ± 4
	[38]	1529 ± 1	188 ± 12	90 ± 3
	[36]	1535 ± 4	166 ± 23	88 ± 6

Table 2. Fit parameters obtained from the fit of the total cross section. The position W_R , the width Γ_R , and the electromagnetic coupling $\sqrt{b_\eta} A_{1/2}^n$ are given for the narrow structure and the $N(1535)$ resonance. The width in the parentheses was extracted from a fit, which has been convoluted with experimental resolution. The indicated errors are statistical.

shown in Fig. 31. The statistical quality is not as good as for the data from [40] but the energy resolution is better. Like the total cross section, they were fit as in [40] with a sum of three Breit-Wigner (BW) functions. A BW function with energy dependent width was used for the $N(1535)$ resonance, a standard BW for the narrow structure, and a much broader BW function was used to describe phenomenologically the remaining contributions from other resonances and non-resonant backgrounds. Due to the simple ansatz, the BW for the $N(1535)$ resonance accounts for the total S_{11} partial wave in this range including contributions from the $N(1650)$ state, some non-resonant backgrounds, the S_{11} - D_{13} interference, and also possible effects from the interference of the $N(1535)$ with the $N(1440)1/2^+$ (Roper resonance). The corresponding integrals of the BW functions are plotted in Fig. 32 as a function of $\cos(\theta_\eta^*)$. The results are quite similar to [40], but better cover the extreme forward and backward angles. As already discussed in [40], the angular dependence of the narrow structure is non-trivial. There is a general trend rising from forward to backward angles. This is what would be expected for a $S_{11} - P_{11}$ interference (the interference term is proportional to $\cos(\theta_\eta^*)$). In comparison to

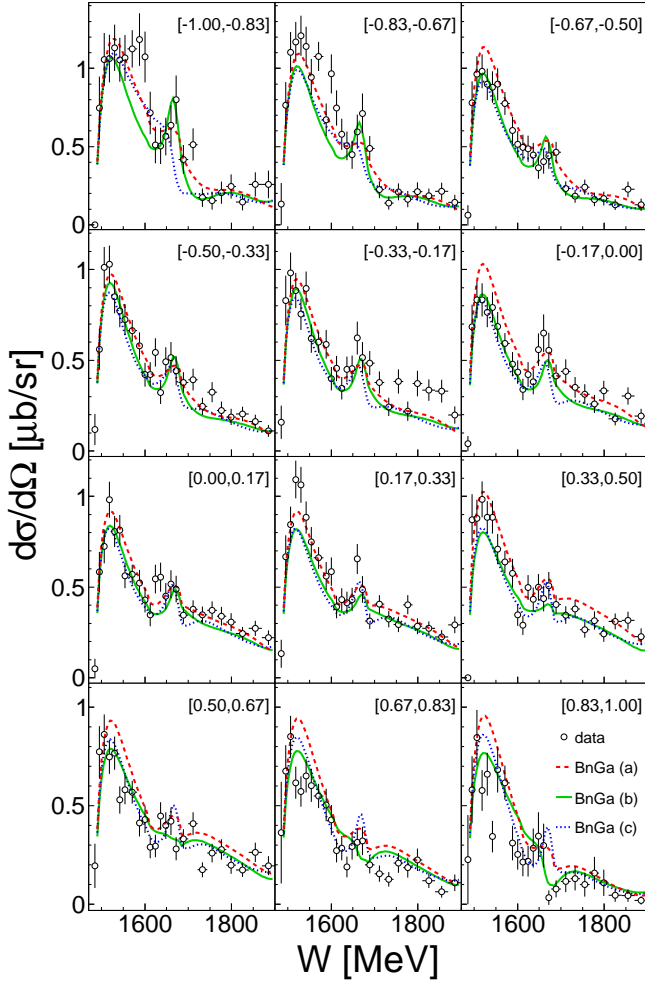


Fig. 33. Differential cross sections in the eta-nucleon cm system for the reaction $\gamma n \rightarrow n\eta$ for 12 angular bins ($\cos(\theta_n^*)$ range indicated in the figure). The present data (black circles) are compared to model calculations from the BnGa group: BnGa (a) (model version with interference in S_{11} wave, dashed red line) [51], BnGa (b) (model with narrow P_{11} resonance with positive $A_{1/2}$ coupling, solid green line) [51], and BnGa (c) (narrow P_{11} resonance with negative $A_{1/2}$ coupling, dotted blue line) [51].

[40], this trend is better visible in the present data due to the better angular coverage. However, like in [40], there is also a maximum around polar angles of 90° degrees, which must have a different origin. This will need further investigation with reaction models such as the BnGa analysis.

As a first step, we compare in Fig. 33 the present data to the calculations of the BnGa group [51], which were obtained by a fit of the data from [40]. Three different fits were discussed in [51]. Fit (a) used only nucleon resonances included in the standard version of the BnGa model and tuned the interference pattern in the S_{11} wave (contributions from $N(1535)$, $N(1650)$, and non-resonant backgrounds) to achieve best agreement with the experimental data. It turned out, that reasonable agreement re-

quires a change in sign (with respect to the value quoted by RPP [46,47]) of the electromagnetic $A_{1/2}$ coupling of the $N(1650)$ for the neutron. In the other two solutions (b), (c), an additional narrow P_{11} resonance was introduced by hand in version (b) so that the interference term between this resonance and the S_{11} partial wave had a positive sign and for (c) so that it had a negative sign. The conclusion was that the version (a) without an additional narrow P_{11} gave the best overall χ^2 and an upper limit was quoted for a contribution of a narrow P_{11} state. However, the present results with better resolution show a clear forward - backward asymmetry of the narrow structure. It is more pronounced at backward angles. This is in better agreement with BnGa solution (b) with the positive sign of an $S_{11} - P_{11}$ interference term than with solution (a) without any narrow state and it clearly contradicts solution (c) with the negative sign of the interference. Also the overall χ^2 is slightly better for the BnGa solution (b) than (a) and (c) (2.43 versus 2.85 and 2.87, respectively). Calculating the χ^2 in the range of the narrow structure, i.e. between 1630 and 1690 MeV, yields values of (a) 2.70, (b) 2.09, and (c) 4.93. Hence, the current results favor a $S_{11} - P_{11}$ interference with positive sign.

A more detailed phenomenological analysis can be done by fitting the Legendre series Eq. 13 to the angular distributions. It is well known that also the angular distributions of the $\gamma p \rightarrow p\eta$ reaction change rapidly shape in this energy region. The ratio of the Legendre coefficients A_1/A_0 has an extremely steep slope around $W=1.7$ GeV and crosses zero shortly below this value [25]. This behavior could also be demonstrated over a significant range of Q^2 by electroproduction [69] of η mesons off protons. The most simple explanation is an $S_{11} - P_{11}$ interference when one of the two partial waves passes rapidly through resonance as discussed in [69]. However, even for the proton this is certainly not the full picture, because also higher order terms show a strong energy dependence in this region. As summarized in [25] also A_3/A_0 changes rapidly sign at the same energy (with the inverse slope compared to A_1/A_0) while A_2/A_0 and A_4/A_0 show pronounced minima.

The Legendre coefficients for $\gamma n \rightarrow n\eta$ up to A_3 from the present work are shown in Fig. 34 and are compared to the published results from [40]. The narrow structure at $W = 1670$ MeV (position indicated by the dashed line) is clearly visible in the A_0 coefficient and consistent with the previous results. However, the structure in the current results is slightly more pronounced, which must originate from the better W resolution in the CBELSA/TAPS experiment. Another interesting finding is the dip structure in A_1 around the same energy. This is clearly more pronounced than in the previous unpolarized data [40], but a similar behavior was seen in the recently published results on the helicity dependent cross section $\sigma_{1/2}$ [55]. The model calculation BnGa (b) with the narrow P_{11} resonance with positive $A_{1/2}$ coupling [51] agrees also with such a behavior. This points to the presence of some $S_{11} - P_{11}$ interference, but like for the proton this is probably not the full explanation of the observed structures.

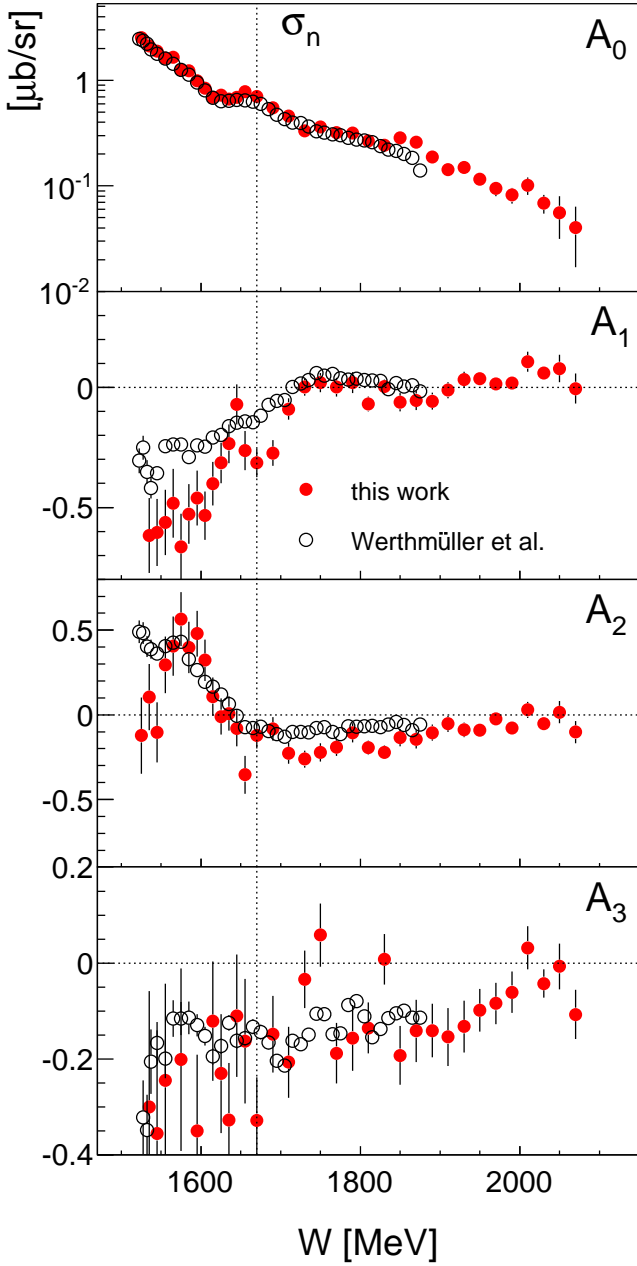


Fig. 34. Legendre coefficients of the angular distributions for the reaction $\gamma n \rightarrow n\eta$. The experimental results from this work (red dots) are compared to previous results from [40] (black circles). The dashed line indicates the position of the narrow structure at $W = 1670$ MeV

4.2 Double Polarization Observable E and Helicity Dependent Cross Sections

The double polarization observable E for the reactions $\gamma p \rightarrow p\eta$ and $\gamma n \rightarrow (n)\eta$ is shown in Fig. 35 as a function of the incident photon energy. Apart from a small systematic deviation (version (1) is slightly higher than version (2)), the two different analysis versions are in good agreement. The experimental results are compared to model predictions from the BnGa analysis [51] and the MAID

model [70]. All models were folded with the Fermi momentum distribution for nucleons bound in deuterium nuclei [66]. The data for the neutron were analyzed in the semi-inclusive way, i.e. all events without detection of any charged particle were accepted as ‘neutron participant’, while for reactions with ‘proton participant’ detection and identification of the recoil proton was required.

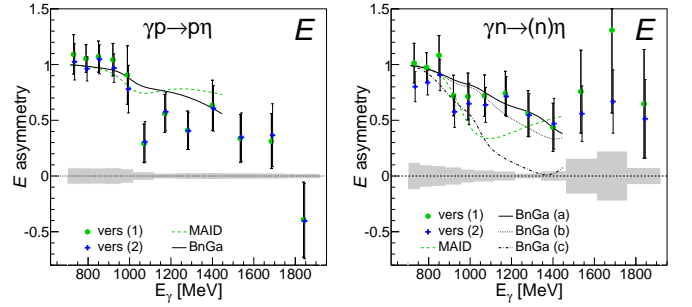


Fig. 35. Double polarization observable E for $\gamma p \rightarrow p\eta$ (left hand side) and $\gamma n \rightarrow (n)\eta$ (right hand side). Green dots: analysis version (1) (carbon subtraction), blue crosses: analysis version (2) (normalization to data from unpolarized deuterium target). Gray shaded areas: systematic uncertainties. Curves: Fermi folded model predictions from MAID (green dashed) [70] and BnGa [51]. For the neutron, three different BnGa calculations are given: BnGa (a) model version with interference in S_{11} wave (solid) [51], BnGa (b) model with narrow P_{11} resonance with positive $A_{1/2}$ coupling (dotted) [51], and BnGa (c) model version with narrow P_{11} resonance with negative $A_{1/2}$ coupling (dashed-dotted) [51]. For better visibility, the data points from version (2) were shifted by +5 MeV with respect to version (1).

Close to threshold, the extracted asymmetries are unity within statistical uncertainties as it is expected from the dominance of the $N(1535)1/2^-$ resonance and predicted by all models. At higher incident photon energies, E drops slightly, indicating contributions from partial waves with $J \geq 3/2$, although contributions from $J = 1/2$ partial waves dominate up to the highest energies. The contribution of $J \geq 3/2$ partial waves above incident photon energies of 1 GeV seems to be more important for the proton than for the neutron, but this is at the edge of statistical significance. The model predictions from BnGa [51] and MAID [70] are similar and in reasonable agreement with the experimental data. For the neutron target, the MAID prediction shows a significant effect from the $N(1675)5/2^-$ resonance at photon energies around 1 GeV, which is not supported by the measurement. The polarization observable E of the BnGa (c) calculation exhibits a stronger fall-off towards higher energies than the other two BnGa models. This behavior is in clear contradiction to the current data.

Even more instructive is the discussion of the helicity dependent cross sections $\sigma_{1/2}$ and $\sigma_{3/2}$, which follow directly from the unpolarized cross section and the E asymmetry. They are summarized in Fig. 36 and the results from the three different analysis methods, which

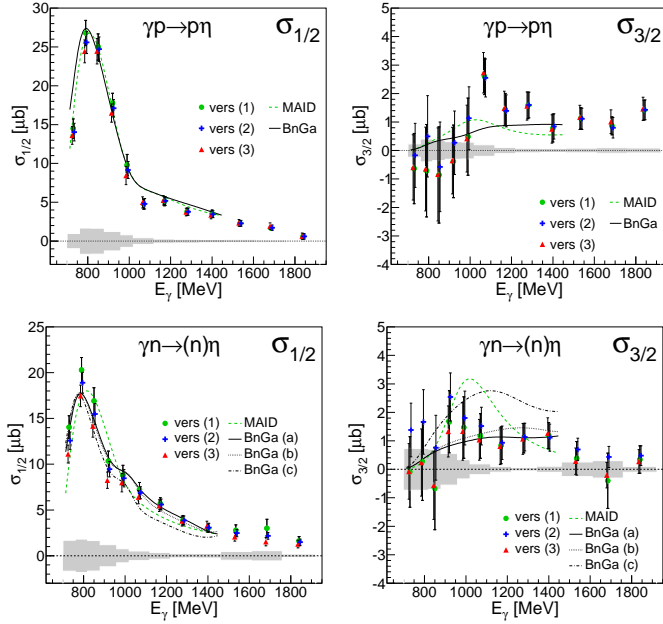


Fig. 36. Helicity dependent cross sections $\sigma_{1/2}$ and $\sigma_{3/2}$ for $\gamma p \rightarrow p\eta$ and $\gamma n \rightarrow (n)\eta$ (reaction and helicity types indicated in figures) from the three different analysis versions (see Sec. 3.7). Gray shaded areas: systematic uncertainties. Curves: Fermi folded model predictions from MAID (green dashed) [70] and BnGa [51]. For the neutron, three different BnGa calculations are given: BnGa (a) model version with interference in S_{11} wave (solid) [51], BnGa (b) model with narrow P_{11} resonance with positive $A_{1/2}$ coupling (dotted) [51], and BnGa (c) model version with narrow P_{11} resonance with negative $A_{1/2}$ coupling (dashed-dotted) [51]. For better visibility, the data points from version (2) and version (3) were shifted by ± 5 MeV with respect to version (1).

agree very well, are compared. The strong dominance of the $\sigma_{1/2}$ contribution for proton and neutron is evident. The structure in the neutron excitation function around 1 GeV appears only in the $\sigma_{1/2}$ part as a bump (broad due to the smearing of the Fermi motion). The $\sigma_{3/2}$ contribution is much smaller in this energy range (on the level of approximately one microbarn) and does not show any peak-like structure as for example predicted by the MAID model due to the $N(1675)5/2^-$ contribution or the BnGa (c) model with the interference of a P_{11} resonance with negative $A_{1/2}$ coupling. Note that all structures are much less narrow than for the above discussed unpolarized cross sections because the Fermi motion was not unfolded from the data. Due to the limited statistics, the neutron data had to be analyzed without a coincident detection of the recoil neutron and effects from Fermi motion could not be removed by an analysis of the final state kinematics. The results agree with the MAMI data published in [55] in so far as the narrow structure appears only in $\sigma_{3/2}$. A direct comparison to the MAMI data is not yet possible because from MAMI only kinematically reconstructed data have been published but not yet the Fermi smeared data.

The proton data shown in Fig. 36 were analyzed in the same way for better comparison. However, for the pro-

ton, the final state kinematics were completely determined because the recoil protons were detected. Therefore, this reaction was also analyzed analogously to the unpolarized cross section by kinematic reconstruction of the final state. The corresponding results as a function of total cm energy W are shown in Fig. 37. They are compared to the original model predictions from MAID [70] and BnGa [51] (without Fermi smearing). The agreement with both predictions for the $\sigma_{1/2}$ contribution is excellent and the very small $\sigma_{3/2}$ contribution is also well described in magnitude. The data for $\sigma_{3/2}$ are within statistical errors not significantly different from zero up to $W \approx 1.65$ GeV. At higher energies, small contributions, for example from the $N(1720)3/2^+$, are possible.

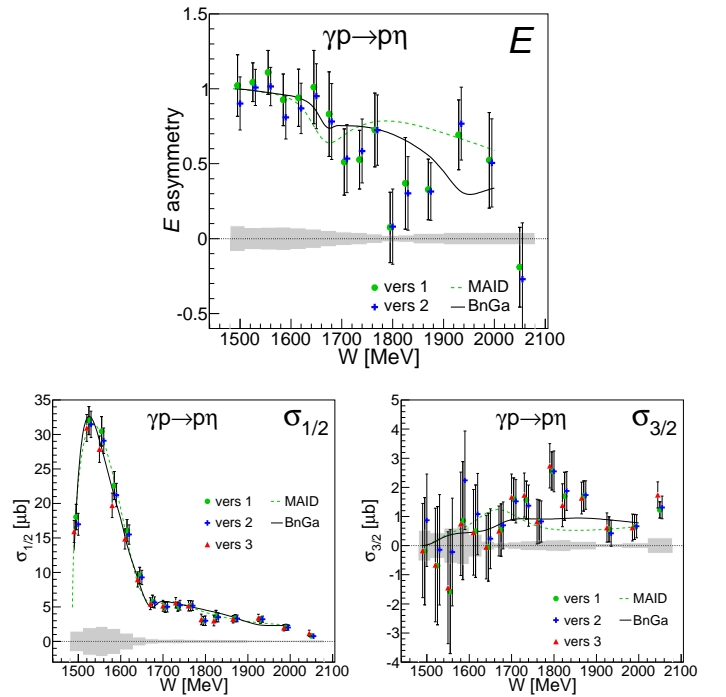


Fig. 37. Double polarization observable E and the helicity dependent cross sections $\sigma_{1/2}$ and $\sigma_{3/2}$ for the reaction $\gamma p \rightarrow p\eta$ as a function of the final state invariant mass. Curves: model predictions from MAID (green dashed) [70] and BnGa (model based on S_{11} interference) [51] for the free proton. For better visibility, the points from version (2) and version (3) were shifted by ± 5 MeV with respect to version (1).

5 Summary and Conclusions

In this paper, total cross sections and angular distributions were presented for η photoproduction from quasifree protons and neutrons. Total cross sections measured as a function of incident photon energy without correction for Fermi smearing agree with previous data [36, 40]. The data after correction for Fermi motion agree much better with [40] than with [36]. The position and width of the narrow

structure seen in previous measurements of the $\gamma n \rightarrow n\eta$ reaction was also confirmed. Angular distributions have shown that the structure is more distinct for backward angles in the cm of the η meson and the neutron and less pronounced in the forward direction. This is a behavior which would agree with an interference between a P_{11} and the strongly dominant S_{11} partial wave. However, the angular dependence is more complicated with a pronounced maximum around 90° . Nevertheless, as a calculation of the χ^2 values showed, the data seem to be in better agreement with the BnGa model solution [51] including an additional narrow P_{11} state than without. This question needs further investigation.

More stringent constraints on the model analyses can be obtained from the measurement of observables exploiting the polarization degrees of freedom of the nucleons and the incident photons. The present work includes results for the double polarization observable E measured with a longitudinally polarized target and a circularly polarized photon beam. This observable, together with the unpolarized cross section σ_0 , allows to split the reaction into contributions with antiparallel ($\sigma_{1/2}$) and parallel ($\sigma_{3/2}$) orientation of nucleon and incident photon spin. The results show that for proton and neutron targets, η photoproduction is strongly dominated by $\sigma_{1/2}$. This was expected for the threshold range, where excitation of the $N(1535)1/2^-$ resonance dominates, but contributions from states with $J \geq 3/2$ also seem to be small at higher incident photon energies. The data for the neutron show that the narrow structure around $E_\gamma \approx 1$ GeV, observed in the unpolarized cross section, is a feature of $\sigma_{1/2}$ and thus very likely related to the S_{11} and/or P_{11} partial waves.

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